

Ohm's law and Continuity Equation:

$$\int (\nabla \cdot \mathbf{J}) dV = \oint \mathbf{J} \cdot d\mathbf{s} = I = \frac{dQ}{dt} = \frac{d}{dt} \int \rho_V dV = \int \frac{\partial \rho_V}{\partial t} dV$$

$$\text{i.e. } I = \frac{dQ}{dt} = \frac{d}{dt} \int \rho_V dV = \int \frac{\partial \rho_V}{\partial t} dV$$

$$\Rightarrow I = \oint \mathbf{J} \cdot d\mathbf{s} = \int (\nabla \cdot \mathbf{J}) dV$$

Divergence theorem

$$\boxed{\nabla \cdot \mathbf{J} = \frac{\partial \rho_V}{\partial t}}$$

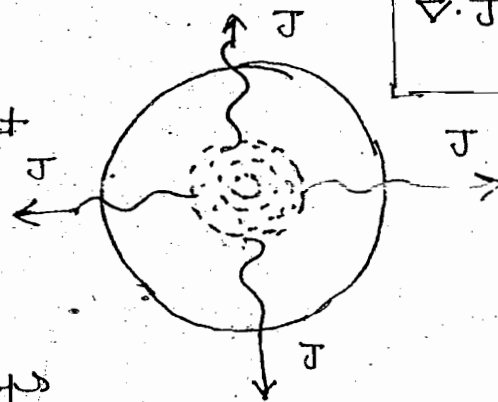
→ Current has outflow and divergence at a rate depending on decrease of volume charges with time

→ Practically

Entering current = Leaving current

charge can neither be created nor be destroyed

$$\left. \begin{aligned} \oint \mathbf{J} \cdot d\mathbf{s} &= 0 \\ \nabla \cdot \mathbf{J} &= 0 \end{aligned} \right\} \rightarrow \text{Always}$$



→ For a uni-directional current flow

$$\frac{\partial J}{\partial x} = \frac{\partial \rho_V}{\partial t}$$

$$\partial J = \partial \rho_V \frac{\partial x}{\partial t}$$

Integrating on both sides

$$J = \rho_V v_d$$

where $v_d = \frac{\partial x}{\partial t} = \text{drift velocity}$

$$V_f = \text{free velocity} = \sqrt{\frac{2qV}{m}}$$

= vacuum tubes, CRO

= few $10^6 - 10^7$ m/s

$$V_f \propto \sqrt{E}$$

where $V_d = \frac{\Delta x}{\Delta t}$ = drift velocity

= change in a material (bulk)

= few cm/sec

$$V_d \propto E \Rightarrow \boxed{V_d = \mu E}$$

↓
mobility of the particle

$$J = \rho_v V_d$$

$$J = \rho_v \mu E$$

$$\Rightarrow \boxed{J = \sigma E} \rightarrow \text{Ohm's law in point form}$$

where σ = conductivity (mho/m) = $\rho_v \mu$

= ability to allow current

= available free carriers $\times$ mobility of carriers

Case-(1) :-

$\sigma = \infty \rightarrow$ very good conductor

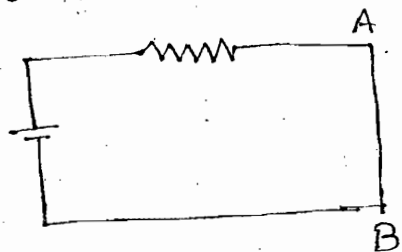
$$\frac{J}{\sigma} = E = 0$$

$\rightarrow$ Electric field cannot exist in a very good conductor

→ $V = \int \mathbf{E} \cdot d\mathbf{l} = \text{constant}$;

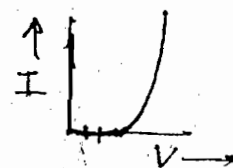
conductor is always a equi-potential region

→ Potential difference or voltage cannot exist in a very good conductor only current can flow but voltage difference good conductor cannot exist



$V_{AB} = 0$ but current flows

$V \uparrow \quad V_{AB} = 0 \quad I_{AB} \uparrow$



eg: → (i) diode current after cutting voltage

(ii) Accumulation — does not exist — only flow exists

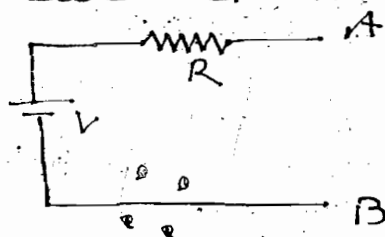
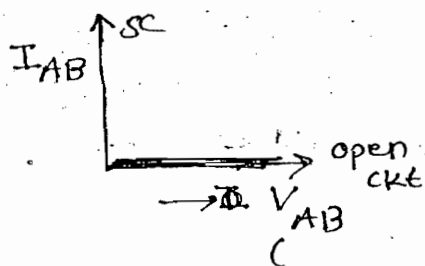
Case — (ii) : —

$\sigma = 0$ very good dielectric.

$J = 0$

→ current cannot exist in a very good dielectric

→ Voltage or potential difference or Electric field can exist but flow current exist



→ Accumulation can exist but flow can't exist.

$$\nabla \cdot \mathbf{J} = - \frac{\partial \rho_v}{\partial t}$$

$$\mathbf{J} = \sigma \mathbf{E}$$

$$\nabla \cdot (\sigma \mathbf{E}) = - \frac{\partial \rho_v}{\partial t}$$

$$\Rightarrow \sigma (\nabla \cdot \mathbf{E}) = - \frac{\partial \rho_v}{\partial t}$$

The ρ_v in time solution is

$$\rho_v(t) = \rho_{v0} \cdot e^{-\frac{t}{\epsilon}}$$

Note:-

Every charge density exponentially spread on the medium at a rate depending on $\frac{\epsilon}{\sigma}$ that called the relaxation time

$$\frac{\epsilon}{\sigma} = \text{Relaxation time} = \frac{\text{Farad}}{\text{m} \times \frac{\text{mho}}{\text{m}}} = \text{ohms} \times \text{farad}$$

Boundary Conditions at conductor surfaces = Second: —

(i) $E_{t1} = E_{t2}$ (General)

As $\sigma = \infty$ along the conductor

$$\frac{J}{\sigma} = E = 0 \text{ along the surface}$$

$$\Rightarrow \boxed{E_{\text{tang.}} = 0}$$

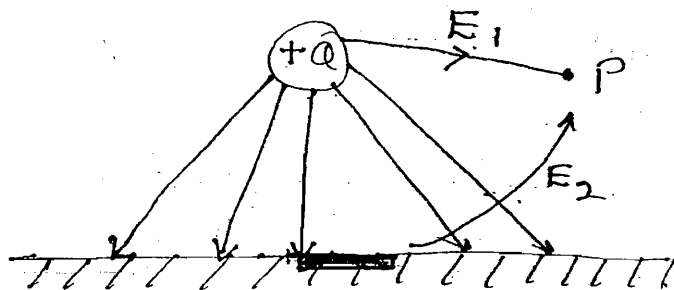
(ii) $\rho_{n2} - \rho_{n1} = \rho_s$ (General)

Electric field can be normal to a conductor which depends on ρ_s on the surface

$$\boxed{E_{\text{normal}} = \rho_s}$$

Conductors, Induction, Method of Images:-

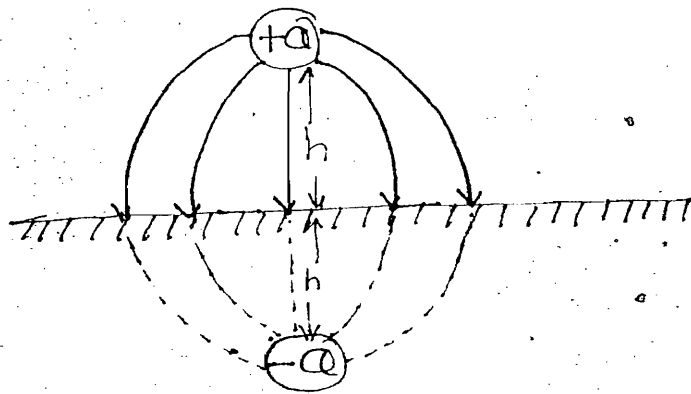
→ When a charge is placed near a conductor surface the field is as shown below



→ The field has tangential component which displaces the free carrier. Hence they are periodically accumulated all ~~the other~~ along the conductor surface. This is called as induced charge.

→ The resultant field at any point is the vector sum of field due to actual charge and induced charge.

→ This field is such that the tangential components are removed and has only normal components as shown below.



→ The field appears to be a dipole field with negative charge called as image charge below the conductor surface.

Summary:-

Every charge and its induction effects are represented by a image charge below the conductor obeying all rules of light and optics.

44.
$$\vec{D}_1 = 2(a_x - \sqrt{3}a_z) \text{ C/m}^2 = \rho_s$$

$$|\rho_s| = |\vec{D}_n| = 2\sqrt{1+3} = 4$$

45.
$$E_n = 2 \text{ V/m}$$

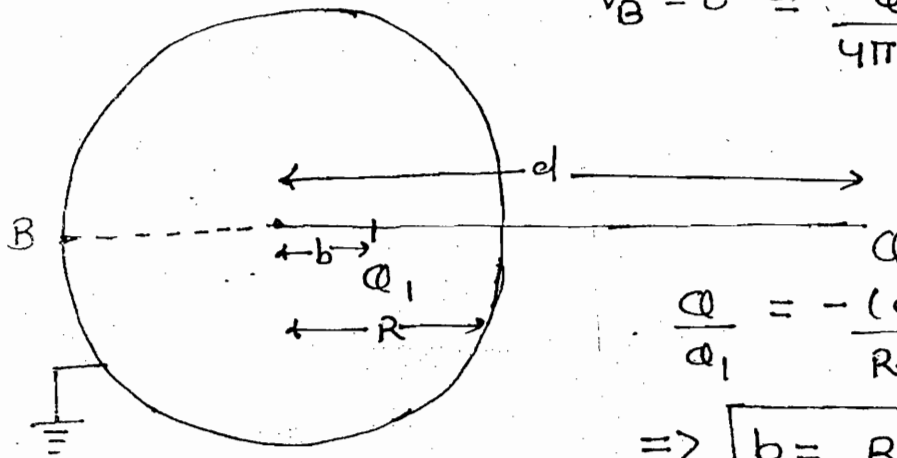
$$\vec{D}_n = \epsilon_0 \epsilon_0 E_n = \rho_s$$

$$\Rightarrow \rho_s = 80 \times 8.8 \times 10^{-12} \times 2$$

$$= 1.41 \times 10^{-9} \text{ C/m}^2$$

$$V_A = 0 = \frac{Q}{4\pi\epsilon(d-R)} + \frac{Q_1}{4\pi\epsilon(R-b)}$$

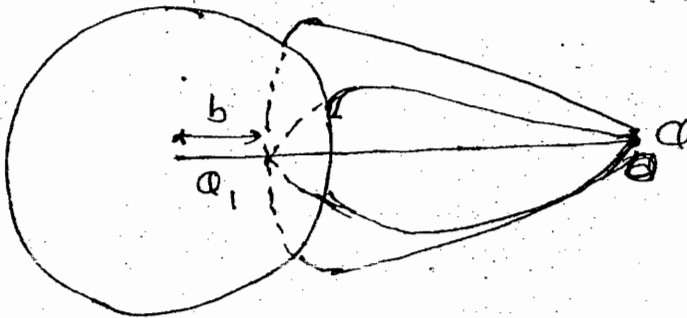
$$V_B = 0 = \frac{Q}{4\pi\epsilon(d+R)} + \frac{Q_1}{4\pi\epsilon(R+b)}$$



$$\frac{Q}{Q_1} = -\frac{(d-R)}{R-b} = -\frac{(d+R)}{(R+b)}$$

$$\Rightarrow \boxed{b = \frac{R^2}{d}} \quad \& \quad \boxed{Q_1 = -\frac{R}{d}Q}$$

Note:-



7. $4\mu C$, Ans

Energy Density in Electric fields ($\frac{dW_E}{dv} = \frac{1}{2} \epsilon E^2$) :-

Consider $Q_1, Q_2, Q_3, \dots, Q_n$ point charges assemble in a region

Total energy of the Electric field = Total Energy expended in assembling the charges

$$W_1 = 0$$

$$W_2 = -Q_2 V_{21}$$

$$W_3 = -Q_3 V_{31} - Q_3 V_{32}$$

$$W_4 = -Q_4 V_{41} - Q_4 V_{42} - Q_4 V_{43}$$

$\vdots$

$$W_n = -Q_n V_{n1} - Q_n V_{n2} - \dots - Q_n V_{nn-1}$$

V_{21} = Potential at 2nd due to 1st charge

$$Q_2 V_{21} = Q_2 \cdot \frac{Q_1}{4\pi\epsilon V_{21}} = Q_1 V_{12}$$

Substitute subscripts can be interchange without change in meaning of value

$$W_1 = 0$$

$$W_2 = -Q_1 V_2$$

$$W_3 = -Q_1 V_{13} - Q_2 V_{23}$$

$$W_4 = -Q_1 V_{14} - Q_2 V_{24} - Q_3 V_{34}$$

⋮

$$W_n = -Q_1 V_{1n} - Q_2 V_{2n} - \dots - Q_{n-1} V_{n-1,n}$$

$$W_E = W_1 + W_2 + W_3 + \dots + W_n$$

total energy

$$2W_E = -Q_1 V_1 - Q_2 V_2 - Q_3 V_3 - \dots - Q_n V_n$$

$$\Rightarrow W_E = -\frac{1}{2} \sum_{i=1}^n Q_i V_i$$

For a continuous charge distribution

$$W_E = -\frac{1}{2} \int \rho_V V dV = -\frac{1}{2} \int (\nabla \cdot \mathbf{D}) V dV = \int \frac{1}{2} \mathbf{D} \cdot (-\nabla V) dV$$

$$\Rightarrow W_E = \int \frac{1}{2} (\mathbf{D} \cdot \mathbf{E}) dV$$

$$\frac{dW_E}{dV} = \text{Energy density}$$

$$= \text{Strength of energy at any point}$$

$$= \frac{1}{2} (\mathbf{D} \cdot \mathbf{E}) = \frac{1}{2} \epsilon E^2$$

Extension:-

$$\frac{dW_H}{dV} = \text{Magnetic Energy density}$$

$$= \frac{1}{2} \mathbf{B} \cdot \mathbf{H} = \frac{1}{2} \mu H^2$$

Capacitors and Inductors:-

→ Capacitance is the ability to confine Electric field in a finitely small region:

$$C = \text{Farad} = \frac{\oint \mathbf{D} \cdot \mathbf{d}\mathbf{s}}{\int \mathbf{E} \cdot \mathbf{d}\mathbf{l}} = \frac{\epsilon \oint \mathbf{E} \cdot \mathbf{d}\mathbf{s}}{\int \mathbf{E} \cdot \mathbf{d}\mathbf{l}} = \frac{Q}{V}$$

→ It is the ratio with charge utilized to the potential developed by the charge.

Specific Geometries:-

- Parallel Plates
- Cocentric cylinders
- Cocentric sphere

Inductance:-

Inductance is the ability to confine H field in a finitely small region

$$L = \text{Henry} = \frac{\int \mathbf{B} \cdot \mathbf{d}\mathbf{s}}{\oint \mathbf{H} \cdot \mathbf{d}\mathbf{l}} = \frac{\mu \int \mathbf{H} \cdot \mathbf{d}\mathbf{s}}{\oint \mathbf{H} \cdot \mathbf{d}\mathbf{l}} = \frac{\Psi_m}{I}$$

→ It is the ratio of the flux developed to the current utilized by the flux

Specific Geometries:-

- Solenoids
- Cocentric cylinder
- Toroids

Parallel Plate Capacitors:-

Parallel Plate Capacitors:-

$$E = \frac{\rho_s}{\epsilon}$$

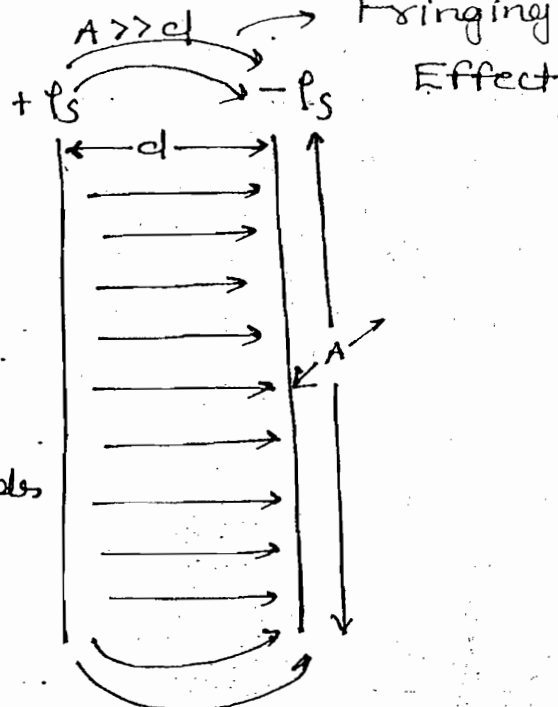
$$C = \frac{Q}{V} = \frac{\rho_s A}{\frac{\rho_s d}{\epsilon}} = \frac{\epsilon A}{d}$$

Capacitance is independent of Q or V and always depends on Area and distance or physical dimensions

$$W_E = \frac{1}{2} \epsilon E^2 (Ad)$$

$$= \frac{1}{2} \frac{\epsilon A}{d} (Ed)^2$$

$$= \frac{1}{2} CV^2 = \frac{1}{2} \frac{Q^2}{C} = \frac{1}{2} QV$$



Multiple Dielectrics in Capacitors:-

Case-(1):-

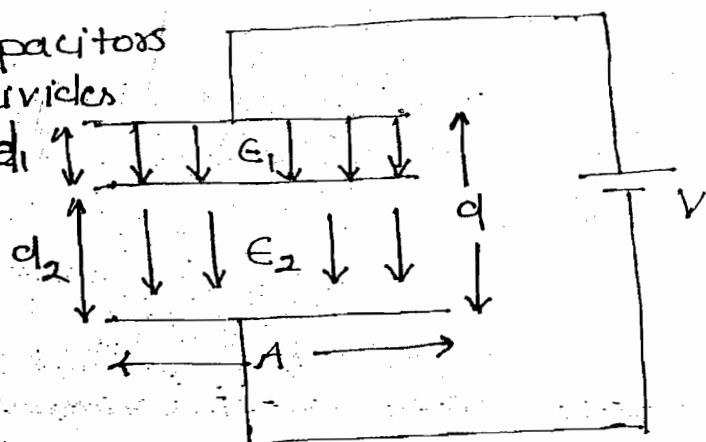
Equal areas of cross-section unequal width of the dielectrics:-

They are two capacitors in series as voltage dividers

$$C_{eq} = \frac{C_1 C_2}{C_1 + C_2}$$

$$= \frac{\epsilon_1 A}{d_1} \cdot \frac{\epsilon_2 A}{d_2}$$

$$\frac{\epsilon_1 A}{d_1} + \frac{\epsilon_2 A}{d_2}$$



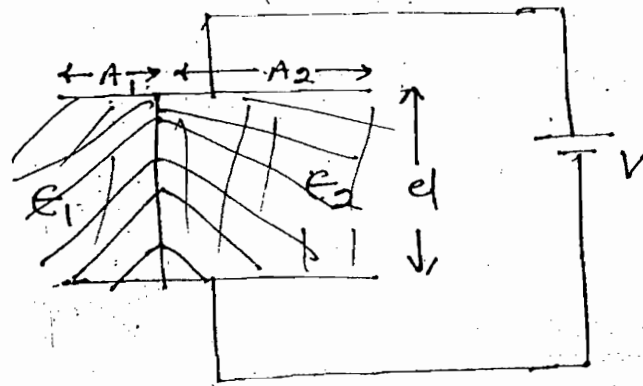
$$V_1 = V_2 \Rightarrow \epsilon_1 E_1 = \epsilon_2 E_2 \Rightarrow \frac{\epsilon_1 V_1}{d_1} = \frac{\epsilon_2 V_2}{d_2}$$

$$\Rightarrow \boxed{\frac{V_1}{V_2} = \left(\frac{\epsilon_2}{\epsilon_1} \right) \left(\frac{d_1}{d_2} \right)}$$

Case-(II):

Unequal areas of cross-section and equal width of the dielectrics :-

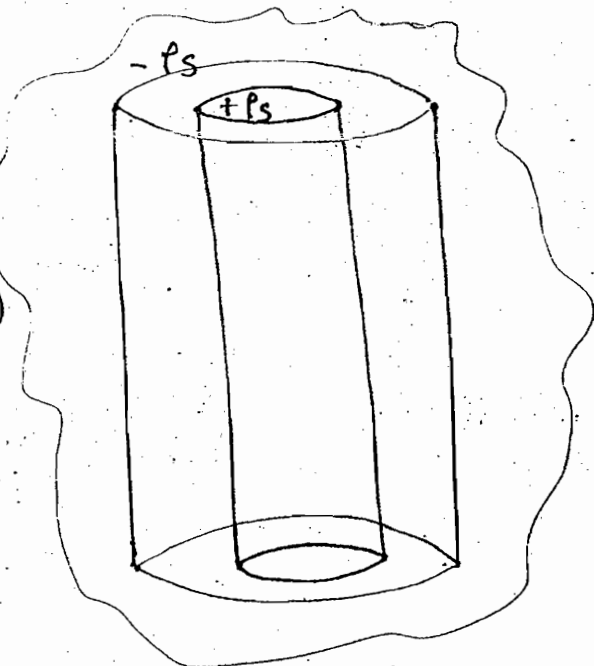
$$C_{eq} = C_1 + C_2$$
$$= \frac{\epsilon_1 A_1 + \epsilon_2 A_2}{d}$$

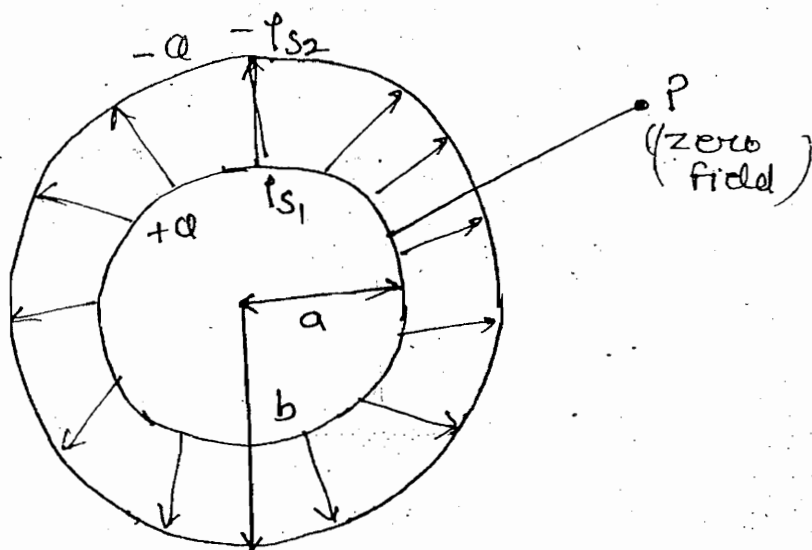


Cocentric Cylinder Capacitance :-

Two cocentric cylinders with equal and opposite charge density may not have confined flux but two equal and opposite charges on concentric cylinders has flux confined only b/w the cylinders

$$Q_{net} = Q_{out} + Q_{in}$$
$$= -\rho_s 2\pi b h + \rho_s 2\pi a h$$
$$\neq 0$$





$$D \propto \rho_s \propto \frac{1}{r}$$

→ The field of a sheet of cylindrical charge obeys same geometry as a line charge

$$C = \frac{Q}{V} = \frac{\rho_L h}{\frac{\rho_L}{2\pi\epsilon} \ln\left(\frac{b}{a}\right)}$$

$$\Rightarrow \boxed{C = \frac{2\pi\epsilon h}{\ln(b/a)}}$$

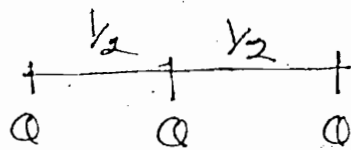
Extension! -

Cocentric sphere capacitance

$$C = \frac{Q}{V}$$

$$\boxed{C = \frac{Q}{\frac{Q}{4\pi\epsilon} \left(\frac{1}{a} - \frac{1}{b} \right)} = \frac{4\pi\epsilon}{\left(\frac{1}{a} - \frac{1}{b} \right)}}$$

$$W_{E1} = 0 + Q \cdot \frac{Q}{4\pi\epsilon_1 \frac{1}{2}}$$



$$+ Q \cdot \frac{Q}{4\pi\epsilon_1 \frac{1}{2}} + Q \cdot \frac{Q}{4\pi\epsilon_1 \cdot 1} = \frac{5Q^2}{4\pi\epsilon_1}$$

$$W_{E2} = \frac{5Q^2}{8\pi\epsilon_1} = \frac{W_{E1}}{2}$$

Note:-

$$W_E \propto QV \propto \frac{1}{r}$$

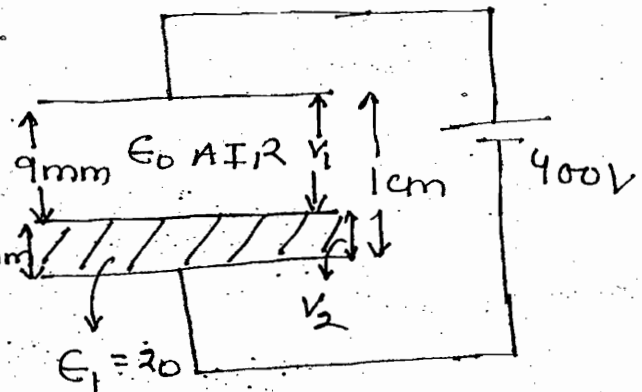
$$\frac{1}{2} (J \cdot A) = \frac{\text{Amp}}{m^2} \times \frac{\text{Joules}}{\text{Amp-m}} = \frac{\text{Joules}}{m^3}$$

$$\Rightarrow \frac{1}{2} (J \cdot A) = \frac{1}{2} ((\nabla \times H) \cdot A) = \frac{1}{2} (H (\nabla \times A)) = \frac{1}{2} B \cdot H$$

$$E_{AIR} = \frac{V_{AIR}}{9mm}$$

$$V_1 + V_2 = 400 \text{ --- (I)}$$

$$\frac{V_1}{V_2} = \left(\frac{\epsilon_2}{\epsilon_1} \right) \left(\frac{d_1}{d_2} \right) \text{ --- (II)}$$



$$\frac{V_1}{V_2} = \frac{20\epsilon_0}{\epsilon_0} \times \frac{9mm}{1mm} = 180$$

$$V_1 = 180V_2$$

$$\text{Ans} \rightarrow V_1 = 44 \text{ kV/m}$$

$$\begin{aligned} \text{Circulation} &= \oint H \cdot dl = I \\ &= I + I + I - (-I) \\ &= 4I, \text{ Ans.} \end{aligned}$$

