

## 4.2 Quantum definitions

### Quantum uncertainty relations

De Broglie relation	$p = \frac{h}{\lambda}$	(4.1)	$p, p$	particle momentum
	$p = \hbar k$	(4.2)	$h$	Planck constant
Planck–Einstein relation	$E = h\nu = \hbar\omega$	(4.3)	$\hbar$	$h/(2\pi)$
Dispersion <sup>a</sup>	$(\Delta a)^2 = \langle(a - \langle a \rangle)^2\rangle$	(4.4)	$\lambda$	de Broglie wavelength
	$= \langle a^2 \rangle - \langle a \rangle^2$	(4.5)	$k$	de Broglie wavevector
General uncertainty relation	$(\Delta a)^2 (\Delta b)^2 \geq \frac{1}{4} \langle \mathbf{i}[\hat{a}, \hat{b}] \rangle^2$	(4.6)	$E$	energy
Momentum–position uncertainty relation <sup>c</sup>	$\Delta p \Delta x \geq \frac{\hbar}{2}$	(4.7)	$\nu$	frequency
Energy–time uncertainty relation	$\Delta E \Delta t \geq \frac{\hbar}{2}$	(4.8)	$\omega$	angular frequency ( $= 2\pi\nu$ )
Number–phase uncertainty relation	$\Delta n \Delta \phi \geq \frac{1}{2}$	(4.9)	$a, b$	observables <sup>b</sup>
			$\langle \cdot \rangle$	expectation value
			$(\Delta a)^2$	dispersion of $a$
			$\hat{a}$	operator for observable $a$
			$[ \cdot, \cdot ]$	commutator (see page 26)
			$x$	particle position
			$t$	time
			$n$	number of photons
			$\phi$	wave phase

<sup>a</sup>Dispersion in quantum physics corresponds to variance in statistics.

<sup>b</sup>An observable is a directly measurable parameter of a system.

<sup>c</sup>Also known as the “Heisenberg uncertainty relation.”

### Wavefunctions

Probability density	$\text{pr}(x, t) dx =  \psi(x, t) ^2 dx$	(4.10)	$\text{pr}$	probability density
Probability density current <sup>a</sup>	$j(x) = \frac{\hbar}{2im} \left( \psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right)$	(4.11)	$\psi$	wavefunction
	$j = \frac{\hbar}{2im} [\psi^*(\mathbf{r}) \nabla \psi(\mathbf{r}) - \psi(\mathbf{r}) \nabla \psi^*(\mathbf{r})]$	(4.12)	$j, j$	probability density current
	$= \frac{1}{m} \Re(\psi^* \hat{\mathbf{p}} \psi)$	(4.13)	$\hbar$	(Planck constant)/(2π)
Continuity equation	$\nabla \cdot \mathbf{j} = -\frac{\partial}{\partial t}(\psi \psi^*)$	(4.14)	$x$	position coordinate
Schrödinger equation	$\hat{H}\psi = i\hbar \frac{\partial \psi}{\partial t}$	(4.15)	$\hat{\mathbf{p}}$	momentum operator
Particle stationary states <sup>b</sup>	$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x)}{\partial x^2} + V(x)\psi(x) = E\psi(x)$	(4.16)	$m$	particle mass
			$\Re$	real part of
			$t$	time
			$H$	Hamiltonian
			$V$	potential energy
			$E$	total energy

<sup>a</sup>For particles. In three dimensions, suitable units would be particles m<sup>-2</sup>s<sup>-1</sup>.

<sup>b</sup>Time-independent Schrödinger equation for a particle, in one dimension.

## Operators

Hermitian conjugate operator	$\int (\hat{a}\phi)^*\psi dx = \int \phi^* \hat{a}\psi dx$	(4.17)	$\hat{a}$ Hermitian conjugate operator $\psi, \phi$ normalisable functions
Position operator	$\hat{x}^n = x^n$	(4.18)	$*$ complex conjugate $x, y$ position coordinates
Momentum operator	$\hat{p}_x^n = \frac{\hbar^n}{i^n} \frac{\partial^n}{\partial x^n}$	(4.19)	$n$ arbitrary integer $\geq 1$ $p_x$ momentum coordinate
Kinetic energy operator	$\hat{T} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2}$	(4.20)	$T$ kinetic energy $\hbar$ (Planck constant)/(2 $\pi$ ) $m$ particle mass
Hamiltonian operator	$\hat{H} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x)$	(4.21)	$H$ Hamiltonian $V$ potential energy
Angular momentum operators	$\hat{L}_z = \hat{x}\hat{p}_y - \hat{y}\hat{p}_x$	(4.22)	$L_z$ angular momentum along $z$ axis (sim. $x$ and $y$ )
	$\hat{L}^2 = \hat{L}_x^2 + \hat{L}_y^2 + \hat{L}_z^2$	(4.23)	$L$ total angular momentum
Parity operator	$\hat{P}\psi(r) = \psi(-r)$	(4.24)	$\hat{P}$ parity operator $r$ position vector

## Expectation value

Expectation value <sup>a</sup>	$\langle a \rangle = \langle \hat{a} \rangle = \int \Psi^* \hat{a} \Psi dx$	(4.25)	$\langle a \rangle$ expectation value of $a$ $\hat{a}$ operator for $a$
	$= \langle \Psi   \hat{a}   \Psi \rangle$	(4.26)	$\Psi$ (spatial) wavefunction $x$ (spatial) coordinate
Time dependence	$\frac{d}{dt} \langle \hat{a} \rangle = \frac{i}{\hbar} \langle [\hat{H}, \hat{a}] \rangle + \left\langle \frac{\partial \hat{a}}{\partial t} \right\rangle$	(4.27)	$t$ time $\hbar$ (Planck constant)/(2 $\pi$ )
Relation to eigenfunctions	if $\hat{a}\psi_n = a_n\psi_n$ and $\Psi = \sum c_n\psi_n$ then $\langle a \rangle = \sum  c_n ^2 a_n$	(4.28)	$\psi_n$ eigenfunctions of $\hat{a}$ $a_n$ eigenvalues $n$ dummy index $c_n$ probability amplitudes
Ehrenfest's theorem	$m \frac{d}{dt} \langle r \rangle = \langle p \rangle$	(4.29)	$m$ particle mass $r$ position vector
	$\frac{d}{dt} \langle p \rangle = -\langle \nabla V \rangle$	(4.30)	$p$ momentum $V$ potential energy

<sup>a</sup>Equation (4.26) uses the Dirac “bra-ket” notation for integrals involving operators. The presence of vertical bars distinguishes this use of angled brackets from that on the left-hand side of the equations. Note that  $\langle a \rangle$  and  $\langle \hat{a} \rangle$  are taken as equivalent.

## Dirac notation

		$n, m$	eigenvector indices
Matrix element <sup>a</sup>	$a_{nm} = \int \psi_n^* \hat{a} \psi_m dx \quad (4.31)$	$a_{nm}$	matrix element
	$= \langle n   \hat{a}   m \rangle \quad (4.32)$	$\psi_n$	basis states
		$\hat{a}$	operator
		$x$	spatial coordinate
Bra vector	bra state vector = $\langle n  $	$\langle \cdot  $	bra
Ket vector	ket state vector = $ m\rangle$	$ \cdot\rangle$	ket
Scalar product	$\langle n   m \rangle = \int \psi_n^* \psi_m dx \quad (4.35)$		
Expectation	if $\Psi = \sum_n c_n \psi_n$	$\Psi$	wavefunction
	then $\langle a \rangle = \sum_m \sum_n c_n^* c_m a_{nm} \quad (4.37)$	$c_n$	probability amplitudes

<sup>a</sup>The Dirac bracket,  $\langle n | \hat{a} | m \rangle$ , can also be written  $\langle \psi_n | \hat{a} | \psi_m \rangle$ .