

## Lecture - 10

Note:-

$R, L, G_1, C \rightarrow$  Primary constant —  $\frac{V}{I}$  Waves

Transmission  
line

$\mu, \sigma, \epsilon \rightarrow$  Material constant —  $E/H$  Waves

Materials

$$\gamma = \sqrt{(R+j\omega L)(G_1+j\omega C)}$$

$$Y = \sqrt{j\omega \mu (\sigma + j\omega \epsilon)}$$

$$Z_0 = \sqrt{\frac{R+j\omega L}{G_1+j\omega C}}$$

$$n = \sqrt{\frac{j\omega \mu}{\sigma + j\omega \epsilon}}$$

Note:- (1)

$$\mu = \frac{|Y_n|}{\omega}$$

$$\sigma = \text{Real}[Y_n]$$

$$\epsilon = \frac{\text{Imag}[Y_n]}{\omega}$$

Note 2:-

$$n = \sqrt{\frac{j\omega \mu}{\sigma + j\omega \epsilon}} = \text{complex} = \frac{E_x}{H_y} = |n| e^{j\theta}$$

$$= \left| \frac{E_x}{H_y} \right| e^{j\theta}$$

Q.L

→ In general material  $E_x$  and  $H_y$  are orthogonal in directions but have a phase shift or time delay which is  $\eta$ 's phase.

Example :-  $E \rightarrow \sin(\omega t - \beta z) a_{ex}$

$$H \rightarrow \sin(\omega t - \beta z - \theta) a_y$$

Note 3:-

$$\gamma = \alpha + j\beta = \sqrt{j\omega\mu(\sigma + j\omega\epsilon)}$$

$$\alpha^2 - \beta^2 + j2\alpha\beta = j\omega\mu\sigma - \omega^2\mu\epsilon$$

$$\alpha = \omega \sqrt{\frac{\mu\epsilon}{2} \left( \sqrt{1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2} - 1 \right)}$$

$$\beta = \omega \sqrt{\frac{\mu\epsilon}{2} \left( \sqrt{1 + \left(\frac{\sigma}{\omega\epsilon}\right)^2} + 1 \right)}$$

### Loss Tangent or Dissipation Factor :-

$\frac{\sigma}{\omega\epsilon}$  decides the propagation aspects  $\alpha$  or  $\beta$  of the EM Waves.

If  $\frac{\sigma}{\omega\epsilon} \gg 1 \Rightarrow$  very good conductor

$\alpha$  is very large  $\Rightarrow$  EM Wave propagation is very tough

→ EM waves cannot easily propagate in good conductors

→ The propagation aspects depends on  $\frac{\sigma}{\omega\epsilon}$  and hence on frequency -  $\omega$  also

e.g:-  $\frac{\sigma}{\omega\epsilon} \gg 1$  For human body upto light frequency

$\frac{\sigma}{\omega\epsilon} \ll 1$  For human body beyond x-ray frequency

eg:-  $\frac{\sigma}{\omega\epsilon} \gg 1$  For earth upto MHz only

$\frac{\sigma}{\omega\epsilon} \ll 1$  for earth beyond microwaves

summary:-

The term  $\frac{\sigma}{\omega\epsilon} = \left| \frac{J_c}{J_d} \right|$  = loss tangent and hence decides the behaviour of the material

$J_c \gg J_d \Rightarrow$  good conductor

$J_d \gg J_c \Rightarrow$  good dielectric

$$\frac{J_c}{J_d} = \frac{-E}{\epsilon \frac{\partial E}{\partial t}} = \frac{-E}{j\omega\epsilon E} = \frac{\sigma}{\omega\epsilon}$$

Case - (III) :-

EM Wave Propagation in very good conductors  
 $(\sigma \gg \omega\epsilon)$

$$\begin{aligned} Y &= \sqrt{j\omega\mu(\sigma + j\omega\epsilon)} \approx \sqrt{j\omega\mu\sigma} \\ &= \sqrt{\omega\mu\sigma} 190^\circ = \sqrt{\omega\mu\sigma} 1+45^\circ \\ &= \sqrt{\frac{\omega\mu\sigma}{2}} + j\sqrt{\frac{\omega\mu\sigma}{2}} \end{aligned}$$

$$\alpha = \beta = \sqrt{\frac{\omega\mu\sigma}{2}}$$

$$\begin{aligned} n &= \sqrt{\frac{j\omega\mu}{\sigma + j\omega\epsilon}} = \sqrt{\frac{j\omega\mu}{\sigma}} = \sqrt{\frac{\omega\mu}{\sigma}} 190^\circ \\ &= \sqrt{\frac{\omega\mu}{\sigma}} 145^\circ = \sqrt{\frac{\omega\mu}{2\sigma}} + j\sqrt{\frac{\omega\mu}{2\sigma}} \end{aligned}$$

$$R = X = \sqrt{\frac{\omega\mu}{2\sigma}}$$

Note (1) :-

→  $\eta$ 's real part ( $R$ ) = cause of attenuation ( $\alpha$ )

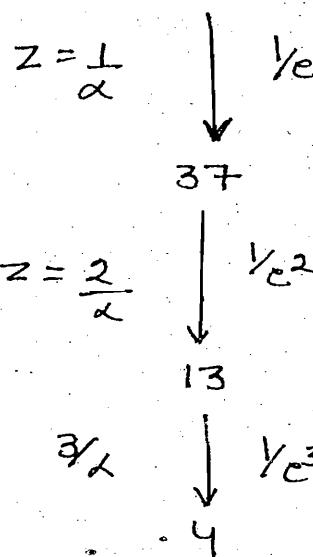
→  $\eta$ 's Imaginary part ( $X$ ) = cause of phase shift ( $\beta$ )

$\eta$ 's phase =  $45^\circ$  means  $E_x$  &  $H_y$  are delayed in time by  $45^\circ$

Skin Depth or Depth of penetration :-

$$|E| = E_0 e^{-\alpha z}$$

$$E_0 = 100$$



→ The first  $\frac{1}{2}$  distance travelled by the wave is where majority of the wave amplitude is attenuated. This distance is called as skin depth. Hence skin depth is the distance travelled by the wave where the wave leakage to  $\frac{1}{e}$  times (37%) of the initial value.

$$S = \frac{1}{\alpha} = \sqrt{\frac{2}{\omega \mu}} = \sqrt{\frac{2\pi}{\omega \mu}} \cdot \frac{1}{\pi} = \frac{1}{\omega R_s}$$

where  $R_s = \eta$ 's real part = skin resistance

$$25. \quad V_p = \frac{\omega}{B} = \frac{\omega}{\lambda} = \omega \delta = 2\pi \times 2 \times 10^5 \times 4 \times 10^{-6}$$

$$= 1.6\pi \text{ m/s} \approx 5 \text{ m/s}$$

$$26. \quad \omega' = 4\omega \quad \delta = \sqrt{\frac{2}{\omega \mu_0}} \Rightarrow \delta' = \frac{\delta}{2}$$

$$V_p' = \omega' \delta' = 4\omega \frac{\delta}{2} = 2V_p = 10 \text{ m/s}$$

Note:-

$$\frac{1}{\omega \epsilon} = 0 \quad \text{Free space} \rightarrow 3 \times 10^8 \text{ m/s}$$

$$\text{Ideal dielectric} \rightarrow \frac{3 \times 10^8}{\sqrt{\epsilon_R}}$$

$$\frac{1}{\omega \epsilon} \uparrow \quad \text{General Material} \rightarrow 10^5 - 10^6 \text{ m/s}$$

$$\frac{1}{\omega \epsilon} \gg 1 \quad \text{Good conductor} \rightarrow \text{few m/s}$$

$$V_p = \frac{\omega}{\sqrt{\frac{\mu \epsilon}{2} \left( \sqrt{1 + \left( \frac{\sigma}{\omega \epsilon} \right)^2} + 1 \right)}}$$

$$\frac{1}{\omega \epsilon} = 1.73$$

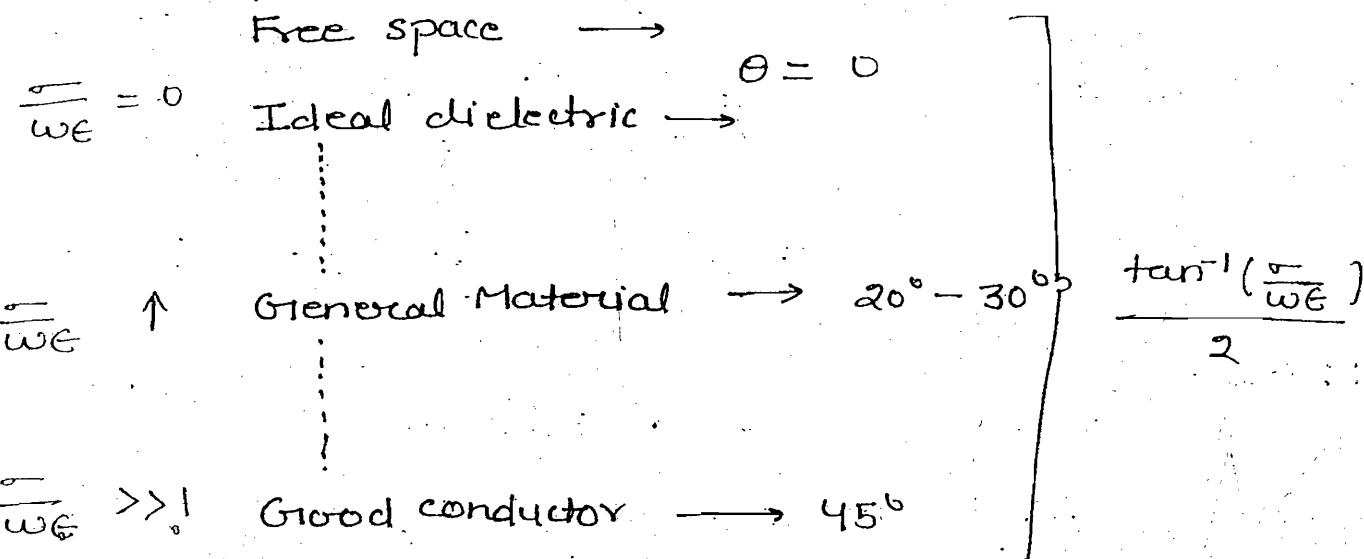
$$\eta' \text{'s phase} = \sqrt{\frac{j\omega \mu}{\sigma + j\omega \epsilon}} = \sqrt{\frac{\omega \mu |90^\circ|}{K \tan^{-1}\left(\frac{\omega \epsilon}{\sigma}\right)}}$$

$$= \sqrt{K |90^\circ - \tan^{-1}\left(\frac{\omega \epsilon}{\sigma}\right)|}$$

$$n' \text{'s phase} = \frac{90^\circ - \tan^{-1}\left(\frac{\omega \epsilon}{\sigma}\right)}{2} = \frac{\tan^{-1}\sqrt{3}}{2} = \frac{60^\circ}{2} = 30^\circ$$

Note:-

$n$ 's phase = Ex to Hy Phase



28.

$$J_c = J_d$$

$$\frac{\sigma}{\omega E} = 1$$

$$\Rightarrow \frac{\sigma}{2\pi f E} = 1$$

$$f = \frac{10^{-2}}{2\pi \times 4 \times 1} = 45 \text{ MHz}$$
  

$$\frac{1}{36\pi \times 10^9}$$

29.

$$\frac{I_c}{I_d} = \frac{J_c}{J_d} = \frac{\sigma}{\omega E}$$

$$\Rightarrow I_d = \frac{I_c \omega E}{58} = \frac{1 \times 2\pi \times 50 \times 8.8 \times 10^{-12}}{58}$$

Ans - (B)

$$n = \sqrt{j\omega \mu} \equiv \sqrt{\omega \mu} 145^\circ$$

$$= R + jX \quad \underline{\text{Ans - 9}}$$

31

$$\eta = 0.02 \quad 145^\circ$$

Ans-(a)

32. Lossy dielectric,  $\frac{\sigma}{\omega\epsilon} \ll 1$

$$\begin{aligned} Y &= \sqrt{j\omega\mu(\sigma + j\omega\epsilon)} = \sqrt{j\omega\mu - j\omega\epsilon \left(1 + \frac{\sigma}{j\omega\epsilon}\right)} \\ &= j\omega\sqrt{\mu\epsilon} \left(1 + \frac{\sigma}{j\omega\epsilon}\right) = \frac{\sigma}{2} \sqrt{\frac{\mu}{\epsilon}} + j\omega\sqrt{\mu\epsilon} \end{aligned}$$

Power, Power density and Poynting Vector :-

→ Energy that can be transformed into other formats is said to be a power format.

$$\frac{\text{Joules}}{\text{second}} = \text{Watts}$$

Using Maxwell's IV Equation

$$\nabla \times H = -\sigma E + \epsilon \frac{\partial E}{\partial t}$$

Take  $\cdot E$  on both sides,

$$(\nabla \times H) \cdot E = -E^2 + \epsilon \frac{\partial E}{\partial t} \cdot E$$

$$\Rightarrow -\nabla \cdot (E \times H) + H \cdot (\nabla \times E) = -E^2 + \epsilon \frac{\partial E}{\partial t} \cdot E$$

$$\Rightarrow -\nabla \cdot (E \times H) - \sigma E^2 = \epsilon \frac{\partial E}{\partial t} \cdot E * \omega \frac{\partial H}{\partial t} \cdot H$$

$$= \frac{1}{\sigma t} \left( \frac{1}{2} \epsilon E^2 + \frac{1}{2} \mu H^2 \right)^2$$

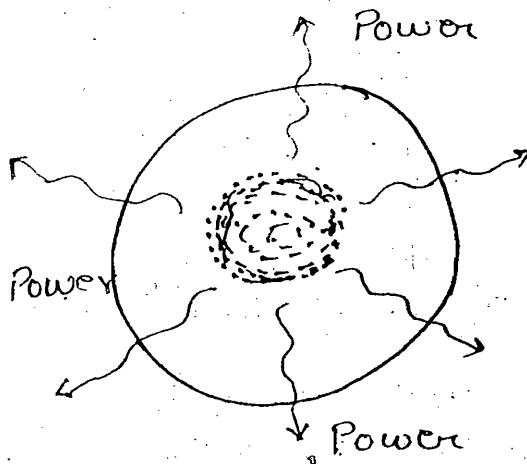
Take  $\int dv$  on both sides

$$\int \nabla \cdot (E \times H) dv + \int -\sigma E^2 dv = -\frac{1}{\sigma t} \int \left( \frac{1}{2} \epsilon E^2 + \frac{1}{2} \mu H^2 \right) dv$$

Using Divergence theorem

$$\oint (E \times H) \cdot dS + \int -E^2 dv = \text{Rate of decrease of } E/H$$

Energy inside a volume



The LHS for a EM wave power crossing the surface

$$\begin{aligned} \text{Hence } E \times H &= \text{Power density of EM Wave} \\ &= \text{Strength of power of EM Wave} \\ &= \text{Poynting Vector of the EM Wave} \\ &= \text{It has the direction of propagation} \end{aligned}$$

$$E(z, t)_x \times H(z, t)_y = P(z, t)_z$$

$$\begin{aligned} &= \text{Power density at a time at a point in space} \\ &= \text{Instantaneous Poynting Vector} \end{aligned}$$

$P(z)_{\text{avg.}}$  = Time Average Poynting vector

$$= E_{\text{RMS}} \cdot H_{\text{RMS}} = \frac{E_0 e^{-\alpha z}}{\sqrt{2}} \cdot \frac{H_0 e^{-\alpha z}}{\sqrt{2}}$$

$$P(z)_{\text{avg.}} = \frac{1}{2} E_0 H_0 e^{-2\alpha z} = \frac{1}{2} \frac{E_0^2}{\eta} e^{-2\alpha z} = \frac{1}{2} \eta H_0^2 e^{-2\alpha z}$$

Note 1:-

The average power decays at  $2\alpha$  rate exponentially in the medium

→ This power lost in the wave is acquired by the medium as ohmic power

$$J \cdot E \text{ or } -E^2 \text{ Watts/m}^3$$

Note 2:-

Alternatively

$$P(z)_{\text{avg.}} = \frac{1}{2} (E \times H^*)$$

\* conjugate operation means average power is independent of harmonic and its phase and depends only on amplitude.

> In free space

$$P_{\text{avg.}} = \frac{1}{2} E_0 H_0 = \frac{1}{2} \frac{E_0^2}{\eta} = \frac{1}{2} \eta H_0^2$$

Workbook! :-

3.  $P_{\text{avg.}} = \frac{1}{2} \eta H_0^2 (-ay)$

$$H_0 = 0.1$$

$$= \frac{1}{2} \cdot 120\pi \cdot (0.1)^2 (-ay) \rightarrow B$$

34.  $W = \int P_{\text{avg.}} ds$   $\curvearrowright x = \text{constant}$

$$= \int \frac{1}{2} \frac{E_0^2}{\eta} a_x \cdot ds \cdot a_x$$

$$= \frac{1}{2} \cdot \frac{60 \cdot 60}{120\pi} \pi \times 4^2 = 240 \text{ W}$$

35. C

36. (1)  $H = 50 \sin(\omega t - \beta z) a_x + 150 \sin(\omega t - \beta z) a_y$

$$P_{\text{avg.}} = \frac{1}{2} \eta H_0^2 (a_z)$$

$$H = 50 \sin(\omega t - \beta z) \left( \frac{a_x + 3a_y}{\sqrt{10}} \right) \sqrt{10}$$

$$= (50\sqrt{10}) \sin(\omega t - \beta z) \left( \frac{a_{0x} + 3a_y}{\sqrt{10}} \right)$$

$$P_{avg} = \frac{1}{2} \cdot 120\pi \cdot (50\sqrt{10})^2 a_z$$

$$H_0 = 50\sqrt{10}$$

$$\begin{aligned} \text{(iii)} \quad H &= 50 \sin(\omega t - \beta z) a_x + 150 \cos(\omega t - \beta z) a_y \\ &= 50 \sin(\omega t - \beta z) a_x (1 + 3j) \\ &= 50 (1 + 3j) \sin(\omega t - \beta z) a_x \end{aligned}$$

Note:-

Orthogonal field components in time <sup>or</sup> space have same effect

$$\text{(iii)} \quad H = 50 \sin(\omega t - \beta z) a_x + 150 \cos(\omega t - \beta z) a_y$$

$$H = 50 \sin(\omega t - \beta z) (a_{0x} + 3j a_y)$$

$$P_{avg} = \frac{1}{2} (E \times H^*)$$

$$E = (50\eta) \sin(\omega t - \beta z) (-a_y) + (150\eta) \cos(\omega t - \beta z) (a_x)$$

$$= (50\eta) \sin(\omega t - \beta z) (3j a_x - a_y)$$

$$P_{avg} = \frac{1}{2} \cdot 50 \cdot 120\pi (3j a_x - a_y) \times 50 (a_x - 3j a_y)$$

$$= \frac{1}{2} \cdot 120\pi \cdot 50^2 (a_z + 3a_z)$$

$$= \frac{1}{2} \cdot 120\pi (50\sqrt{10})^2 a_z$$

$$E = (a_x + j a_y) e^{jkz - j\omega t}$$

$$H = \frac{k}{\omega u} (a_y + j a_x) e^{jkz - j\omega t}$$

$$\begin{aligned} P_{avg.} &= \frac{1}{2} (a_x + j a_y) \times \frac{k}{\omega u} (a_y - j a_x) \\ &= \frac{k}{2\omega u} (a_z - a_z) = 0 \rightarrow \text{Null vector} \end{aligned}$$

Conventional :-

$$E = 100 (j a_x + 2 a_y - j a_z) e^{j\omega t}$$

$$H = (-a_x + j a_y + a_z) e^{j\omega t}$$

$$\begin{aligned} P_{avg.} &= \frac{1}{2} (E \times H^*) = \frac{1}{2} 100 (j a_x + 2 a_y - 2 a_z) \times \\ &\quad (-a_x - j a_y + a_z) \\ &= 50 \begin{vmatrix} a_x & a_y & -a_z \\ j & 2 & -j \\ -1 & -j & 1 \end{vmatrix} \\ &= 50 (3 a_x + 3 a_z) \end{aligned}$$

$$|P_{avg.}| = 150\sqrt{2}$$

$$P_{avg.} \text{ direction} = (a_x + a_z)$$

$$\text{Unit P}_{avg.} \text{ direction} = \frac{a_x + a_z}{\sqrt{2}}$$

## Wave Polarization:-

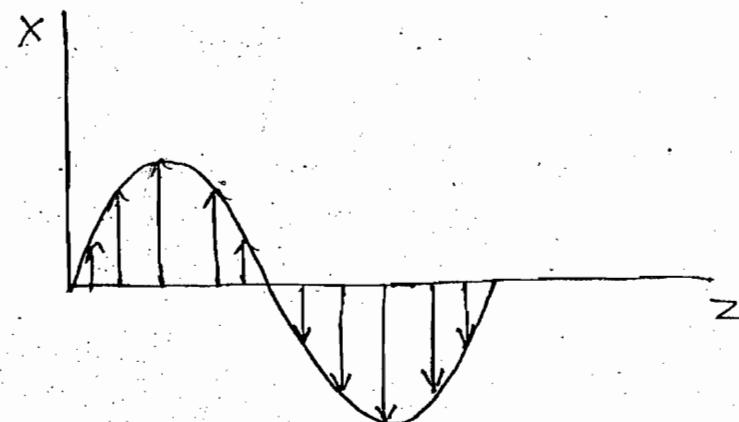
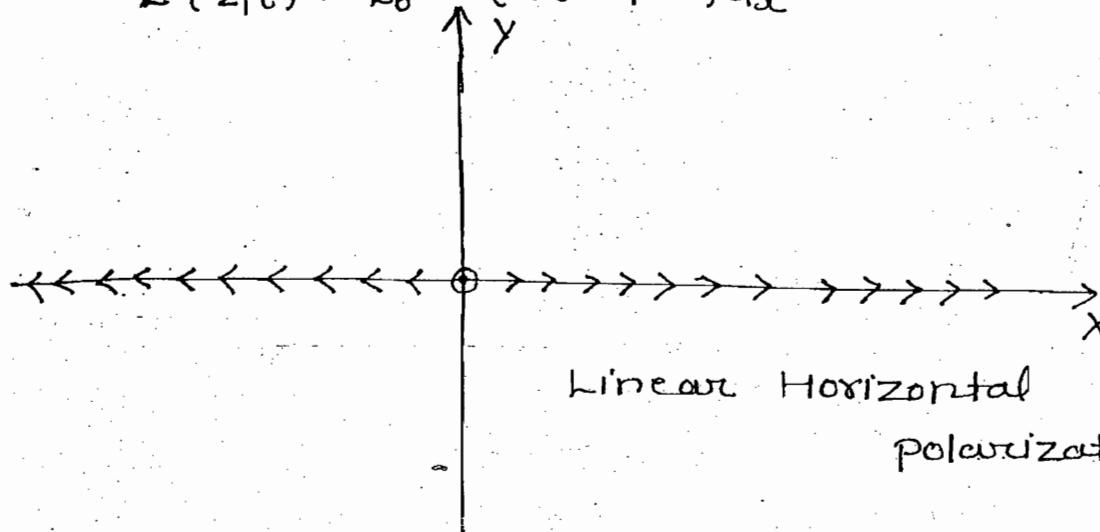
It is the electric field orientation & the possible planar components of the E field satisfying the basic transverse nature

EM Wave is propagating along z-axis in the following cases:-

### Case (I) :-

E field is oriented in x-direction

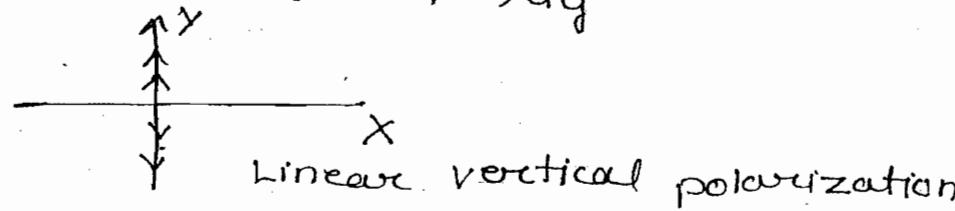
$$E(z,t) = E_0 \sin(\omega t - \beta z) a_x$$



### Case -(II) :-

E field is oriented in y-direction

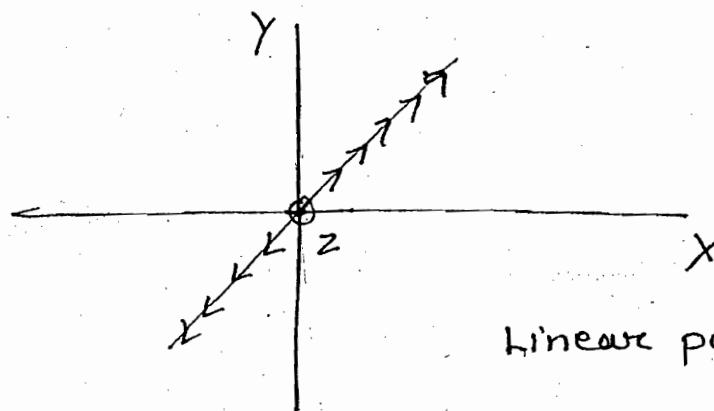
$$E(z,t) = E_0 \sin(\omega t - \beta z) a_y$$



### Case-(III) :-

E field is oriented in x & y directions.

$$E(z,t) = E_1 \sin(\omega t - \beta z) a_x + E_2 \sin(\omega t - \beta z) a_y$$



Linear polarization

### Note:-

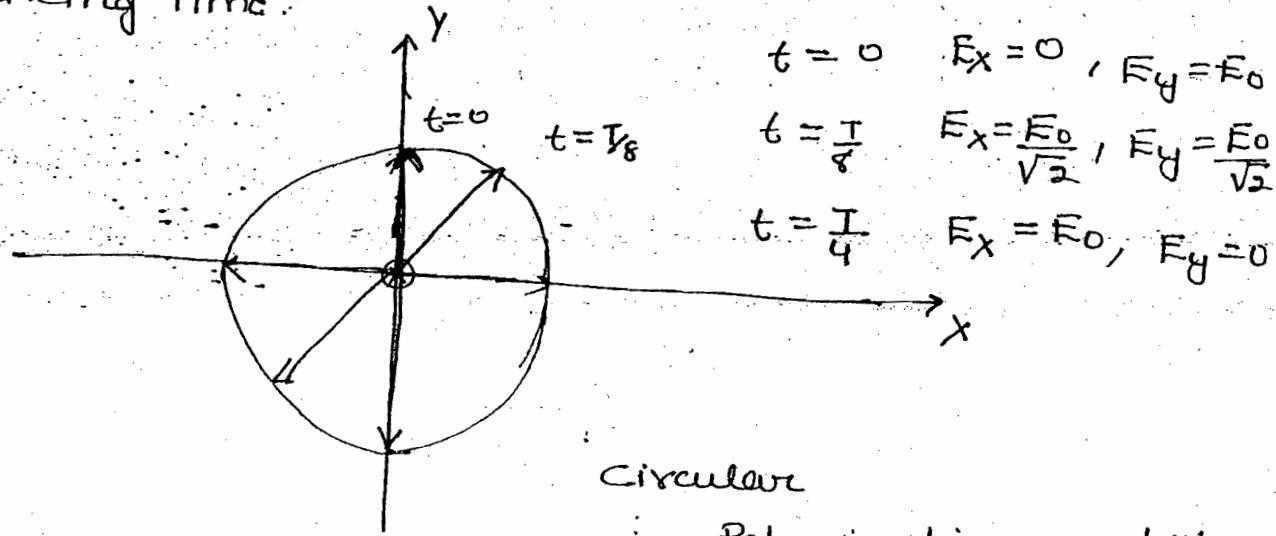
If the EM waves has single E field components or two planar components both inphase the wave is said to be linearly polarized

### Case-(IV) :-

E field is oriented in x & y directions

$$E(z,t) = E_0 \sin(\omega t - \beta z) a_x + E_0 \cos(\omega t - \beta z) a_y$$

→ Polarization is identified by studying the trace of the E field on the  $z=0$  plane for various advancing time.



Circular

Polarization.

## summary:-

If the EM Waves has two planar components of E field both out of phase at  $90^\circ$  and equal amplitude , the wave is circularly polarised.

## Note:-

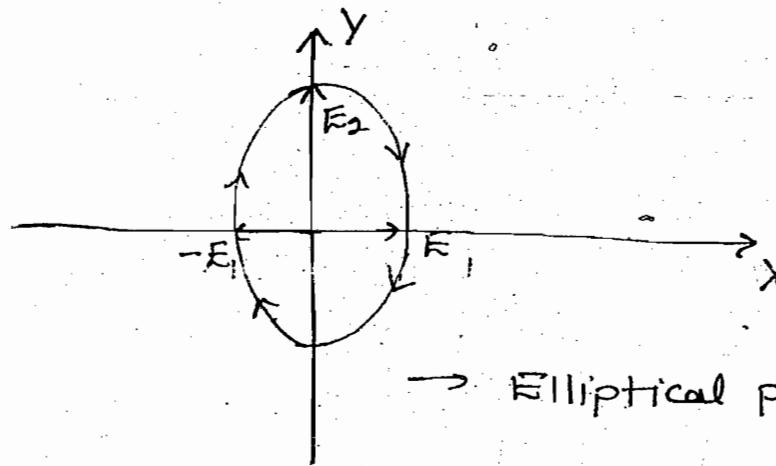
Sense of rotation left or right , if the left hand thumb points towards propagation direction and the closed finger along advancing time, the wave is left circularly polarised

e.g:- out of paper - propagation clockwise time advancement - left circular

## case-(V) :-

E field is oriented in x & y directions

$$\mathbf{E}(z,t) = E_1 \sin(\omega t - \beta z) \mathbf{a}_x + E_2 \cos(\omega t - \beta z) \mathbf{a}_y$$

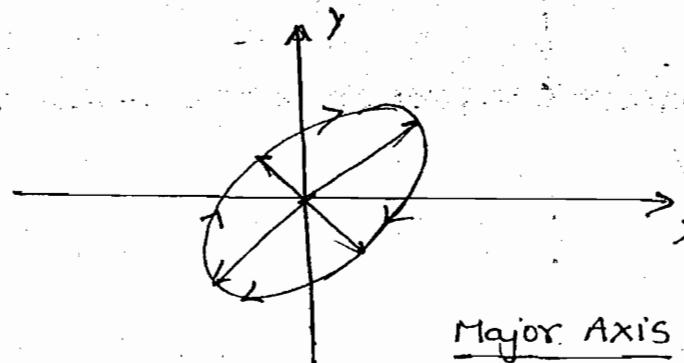


→ Elliptical polarization

## case-(VI) :-

E field is oriented in x & y direction

$$\mathbf{E}(z,t) = E_1 \sin(\omega t - \beta z) \mathbf{a}_x + E_2 \sin(\omega t - \beta z + \theta) \mathbf{a}_y$$



Major Axis = AR = [1, ∞)

Minor Axis

Circle  
Linear

Workbook :-

3 (i)  $\vec{E} = 25 \sin(\omega t + 4x) (a_y + 6a_z)$   
 → Linear

ii)  $\vec{E} = 25 \sin(\omega t + 4x) a_y + 25 \cos(\omega t + 4x) a_z$   
 → Circular

iii)  $\vec{E} = 25 \sin(\omega t + 4x + 60^\circ) a_y + 25 \cos(\omega t + 4x + 60^\circ) a_z$   
 → Circular

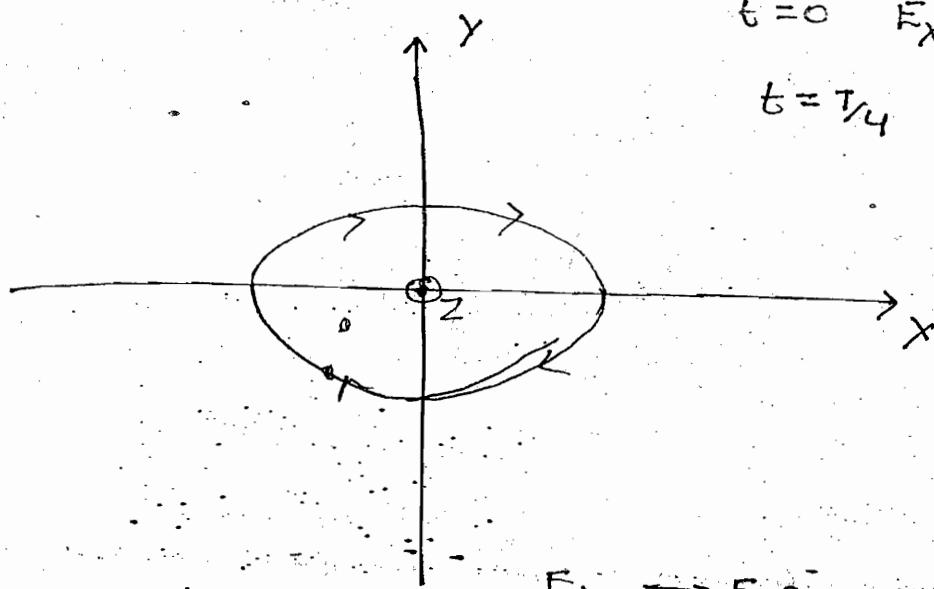
iv)  $\vec{E} = 25 \sin(\omega t + 4x) a_y + (25\sqrt{2}) \sin(\omega t + 4x + 45^\circ) a_z$   
 → Elliptical

v) Elliptical

vi) Circular

g.  $\theta = 60^\circ \rightarrow$  elliptical

o.  $\theta = 90^\circ$



$t=0 \quad E_x = E_1 \neq E_y = 0$

$t=\frac{T}{4} \quad E_x = 0 \quad E_y = -E_1$

Summary:-

$E_x \rightarrow E_1 \cos \omega t \quad E_y \rightarrow E_2 \sin \omega t$

$E_x$	$E_y$	Prop.	Polarization
$\cos$	$-\sin$	$z$	Left
$\cos$	$\sin$	$z$	Right
$\cos$	$-\sin$	$-z$	Right
$-\sin$	$\cos$	$z$	Right

$(ay + jaz) e^{j\theta}$

$\downarrow$   
 $\sin$   
 $\downarrow$   
 $\cos$

$\cos$

$-\sin$

