

4.3 Wave mechanics

Potential step^a

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| Potential function | $V(x) = \begin{cases} 0 & (x < 0) \\ V_0 & (x \geq 0) \end{cases} \quad (4.38)$ |
| Wavenumbers | $\hbar^2 k^2 = 2mE \quad (x < 0) \quad (4.39)$ |
| | $\hbar^2 q^2 = 2m(E - V_0) \quad (x > 0) \quad (4.40)$ |
| Amplitude reflection coefficient | $r = \frac{k - q}{k + q} \quad (4.41)$ |
| Amplitude transmission coefficient | $t = \frac{2k}{k + q} \quad (4.42)$ |
| Probability currents ^b | $j_I = \frac{\hbar k}{m} (1 - r ^2) \quad (4.43)$ |
| | $j_{II} = \frac{\hbar q}{m} t ^2 \quad (4.44)$ |
| V particle potential energy V_0 step height \hbar (Planck constant)/(2 π) k, q particle wavenumbers m particle mass E total particle energy r amplitude reflection coefficient t amplitude transmission coefficient j_I particle flux in zone I j_{II} particle flux in zone II | |

^aOne-dimensional interaction with an incident particle of total energy $E = KE + V$. If $E < V_0$ then q is imaginary and $|r|^2 = 1$. $1/|q|$ is then a measure of the tunnelling depth.

^bParticle flux with the sign of increasing x .

Potential well^a

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| | <p>incident particle →</p> | $V(x)$ I $-a$ II a III x | |
| Potential function | $V(x) = \begin{cases} 0 & (x > a) \\ -V_0 & (x \leq a) \end{cases} \quad (4.45)$ | | V particle potential energy V_0 well depth \hbar (Planck constant)/(2π) $2a$ well width k, q particle wavenumbers m particle mass E total particle energy r amplitude reflection coefficient t amplitude transmission coefficient j_i particle flux in zone I j_{III} particle flux in zone III n integer > 0 E_n Ramsauer energy |
| Wavenumbers | $\hbar^2 k^2 = 2mE \quad (x > a) \quad (4.46)$ | | |
| | $\hbar^2 q^2 = 2m(E + V_0) \quad (x < a) \quad (4.47)$ | | |
| Amplitude reflection coefficient | $r = \frac{i e^{-2ika}}{2kq \cos 2qa - i(q^2 + k^2) \sin 2qa} \quad (4.48)$ | | |
| Amplitude transmission coefficient | $t = \frac{2kqe^{-2ika}}{2kq \cos 2qa - i(q^2 + k^2) \sin 2qa} \quad (4.49)$ | | |
| Probability currents ^b | $j_i = \frac{\hbar k}{m} (1 - r ^2) \quad (4.50)$ | | |
| | $j_{III} = \frac{\hbar k}{m} t ^2 \quad (4.51)$ | | |
| Ramsauer effect ^c | $E_n = -V_0 + \frac{n^2 \hbar^2 \pi^2}{8ma^2} \quad (4.52)$ | | |
| Bound states ($V_0 < E < 0$) ^d | $\tan qa = \begin{cases} k /q & \text{even parity} \\ -q/ k & \text{odd parity} \end{cases} \quad (4.53)$ | | |
| | $q^2 - k ^2 = 2mV_0/\hbar^2 \quad (4.54)$ | | |

^aOne-dimensional interaction with an incident particle of total energy $E = \text{KE} + V > 0$.

^bParticle flux in the sense of increasing x .

^cIncident energy for which $2qa = n\pi$, $|r| = 0$, and $|t| = 1$.

^dWhen $E < 0$, k is purely imaginary. $|k|$ and q are obtained by solving these implicit equations.

Barrier tunnelling^a

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| Potential function | $V(x) = \begin{cases} 0 & (x > a) \\ V_0 & (x \leq a) \end{cases}$ | (4.55) | |
| Wavenumber and tunnelling constant | $\hbar^2 k^2 = 2mE \quad (x > a)$ | (4.56) | |
| | $\hbar^2 \kappa^2 = 2m(V_0 - E) \quad (x < a)$ | (4.57) | |
| Amplitude reflection coefficient | $r = \frac{-ie^{-2ika}(k^2 + \kappa^2)\sinh 2\kappa a}{2\kappa \cosh 2\kappa a - i(k^2 - \kappa^2)\sinh 2\kappa a}$ | (4.58) | |
| Amplitude transmission coefficient | $t = \frac{2\kappa e^{-2ika}}{2\kappa \cosh 2\kappa a - i(k^2 - \kappa^2)\sinh 2\kappa a}$ | (4.59) | |
| Tunnelling probability | $ t ^2 = \frac{4k^2 \kappa^2}{(k^2 + \kappa^2)^2 \sinh^2 2\kappa a + 4k^2 \kappa^2}$ | (4.60) | |
| | $\simeq \frac{16k^2 \kappa^2}{(k^2 + \kappa^2)^2} \exp(-4\kappa a) \quad (t ^2 \ll 1)$ | (4.61) | |
| Probability currents ^b | $j_I = \frac{\hbar k}{m}(1 - r ^2)$ | (4.62) | |
| | $j_{III} = \frac{\hbar k}{m} t ^2$ | (4.63) | |
| | | V particle potential energy V_0 well depth \hbar (Planck constant)/(2π) $2a$ barrier width k incident wavenumber κ tunnelling constant m particle mass E total energy ($< V_0$) r amplitude reflection coefficient t amplitude transmission coefficient $ t ^2$ tunnelling probability j_I particle flux in zone I j_{III} particle flux in zone III | |

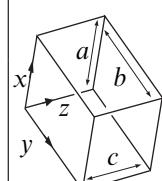
^aBy a particle of total energy $E = KE + V$, through a one-dimensional rectangular potential barrier height $V_0 > E$.

^bParticle flux in the sense of increasing x .

Particle in a rectangular box^a

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| Eigenfunctions | $\Psi_{lmn} = \left(\frac{8}{abc} \right)^{1/2} \sin \frac{l\pi x}{a} \sin \frac{m\pi y}{b} \sin \frac{n\pi z}{c}$ | Ψ_{lmn} eigenfunctions a, b, c box dimensions l, m, n integers ≥ 1 |
| Energy levels | $E_{lmn} = \frac{\hbar^2}{8M} \left(\frac{l^2}{a^2} + \frac{m^2}{b^2} + \frac{n^2}{c^2} \right)$ | E_{lmn} energy \hbar Planck constant M particle mass |
| Density of states | $\rho(E) dE = \frac{4\pi}{\hbar^3} (2M^3 E)^{1/2} dE$ | $\rho(E)$ density of states (per unit volume) |

^aSpinless particle in a rectangular box bounded by the planes $x=0, y=0, z=0, x=a, y=b$, and $z=c$. The potential is zero inside and infinite outside the box.



Harmonic oscillator

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| Schrödinger equation | $-\frac{\hbar^2}{2m} \frac{\partial^2 \psi_n}{\partial x^2} + \frac{1}{2} m \omega^2 x^2 \psi_n = E_n \psi_n \quad (4.67)$ | \hbar (Planck constant)/(2π) m mass ψ_n n th eigenfunction x displacement n integer ≥ 0 ω angular frequency E_n total energy in n th state |
| Energy levels ^a | $E_n = \left(n + \frac{1}{2} \right) \hbar \omega \quad (4.68)$ | H_n Hermite polynomials |
| Eigen-functions | $\psi_n = \frac{H_n(x/a) \exp[-x^2/(2a^2)]}{(n! 2^n a \pi^{1/2})^{1/2}} \quad (4.69)$ where $a = \left(\frac{\hbar}{m\omega} \right)^{1/2}$ | y dummy variable |
| Hermite polynomials | $H_0(y) = 1, \quad H_1(y) = 2y, \quad H_2(y) = 4y^2 - 2$ $H_{n+1}(y) = 2yH_n(y) - 2nH_{n-1}(y) \quad (4.70)$ | |

^a E_0 is the zero-point energy of the oscillator.