

Lecture - 6

Closed line integral of electric field = Maxwell's II Eqn:-

$$\oint \mathbf{E} \cdot d\mathbf{l} = 0$$

- Potential difference exists b/w 2 different points but b/w the same point potential difference is zero
- Potential is unique at a point at a time
- Workdone in moving a charge in any closed path is always zero
- In any closed path energy acquired = energy lost
i.e. Energy acquired in field direction = energy lost opposite to field direction.
- Electric field is a conservative field & never forms closed loops.

i.e. Irrotational vector ($\nabla \times \mathbf{E} = 0$)

- Workdone in moving a charge b/w two points is independent of path of consideration.
- This is KVL in electric fields.

Potential, Vector Potential and Ampere's Law in H field:-

- Potential is called as MMF or Magneto Motive force and is a scalar measure of magnetic field strength at any point in the field.

$$V_m = \text{MMF} = \int \vec{H} \cdot d\vec{l} = \text{Amp}$$

$$\mathbf{H} = \nabla V_m$$

$$\boxed{\oint \mathbf{H} \cdot d\mathbf{l} = I}$$

This is Ampere's law in Integral form

$$\nabla \times H = J$$

This is Ampere's law in point form

Statement:- effects

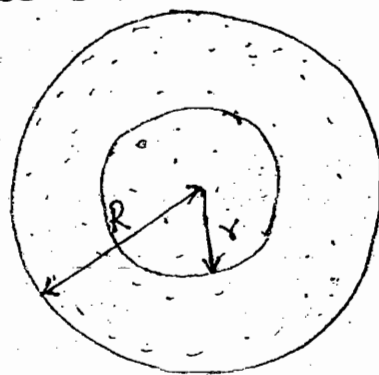
The net circulation of magnetic field in any closed line is always equal to the current crossing the surface enclosed.

$$\begin{aligned} \text{Circulation} &= \text{Strength} \times \text{length} \\ &= \text{current} \end{aligned}$$

consider a closed ~~ring~~ line symmetrical & concentric with the I flow direction such that $H = \text{constant everywhere}$

$$\oint H \cdot dl = I$$

↓
constant



$$H \cdot \text{Length} = \text{current crossing } r < R$$

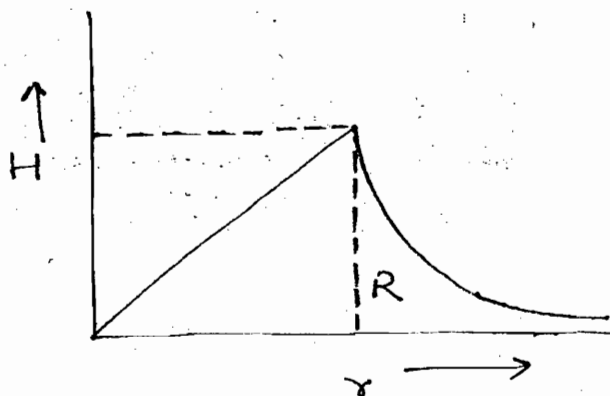
$$H = \frac{\frac{I}{\pi R^2} \cdot \pi r^2}{2\pi r} = \frac{I r}{2\pi R^2}$$

$$r > R$$

$$H = \frac{I}{2\pi r}$$

Note:-

- Point $\rightarrow 1/r^2$
- Line $\rightarrow 1/r$
- Sheet $\rightarrow \text{Uniform}$
- Solid $\rightarrow r$



Vector Potential :-

Limitation of Scalar Potential (MMF) :-

→ The definition of potential in electric field is consistent with its nature i.e. irrotational nature

$$E = -\nabla V$$

$$\nabla \times E = 0$$

$$\nabla \times (-\nabla V) = 0$$

$$\nabla \times (\text{Grad of Scalar}) = 0$$

→ A similar definition in magnetic field for MMF is not consistent with Ampere's law

$$H = \nabla V_m$$

$$\nabla \times H = J$$

$$\nabla \times (\nabla V_m) = 0 = J$$

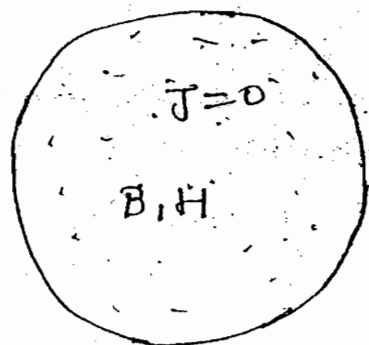
Hence MMF is defined and exists in only those regions where $J=0$

i.e. outside current flowing regions

i.e. free space conditions and outside conductor

eg:- (i) MMF b/w windings of solenoids

(ii) MMF in air gaps in machines



B, H, J

Vector Potential $\vec{A}$:-

→ The basic definition of vector potential is in accordance to nature of magnetic field i.e. solenoidal nature. Hence if $\text{curl of } \vec{A} = \vec{B}$ then $\vec{A}$ is called as vector potential

$$\vec{B} = \nabla \times \vec{A}$$

$$\nabla \cdot \vec{B} = 0$$

$$\nabla \cdot (\nabla \times \vec{A}) = 0$$

$$\text{Div. (curl of vector)} = 0$$

→ The term vector potential has a units weber/m which physically signifies the work for current element

$$\vec{A} = \text{Weber/m} = \frac{\text{Joules}}{\text{Amp-m}}$$

$$= \frac{W}{I d\vec{l}}$$

and as $I d\vec{l}$ is a vector, $\vec{A}$ is a vector quantity.

Note:-

$$\frac{\text{Weber}}{\text{second}} = \text{Volts}$$

$$\frac{\text{Weber}}{\text{second}} = \frac{\text{Joules}}{\text{Coulomb}}$$

$$\frac{\text{Weber}}{\text{m}} = \frac{\text{Joules}}{\text{Amp-m}}$$

$$V = \frac{W}{Q} = \frac{Q}{4\pi\epsilon_0 r}$$

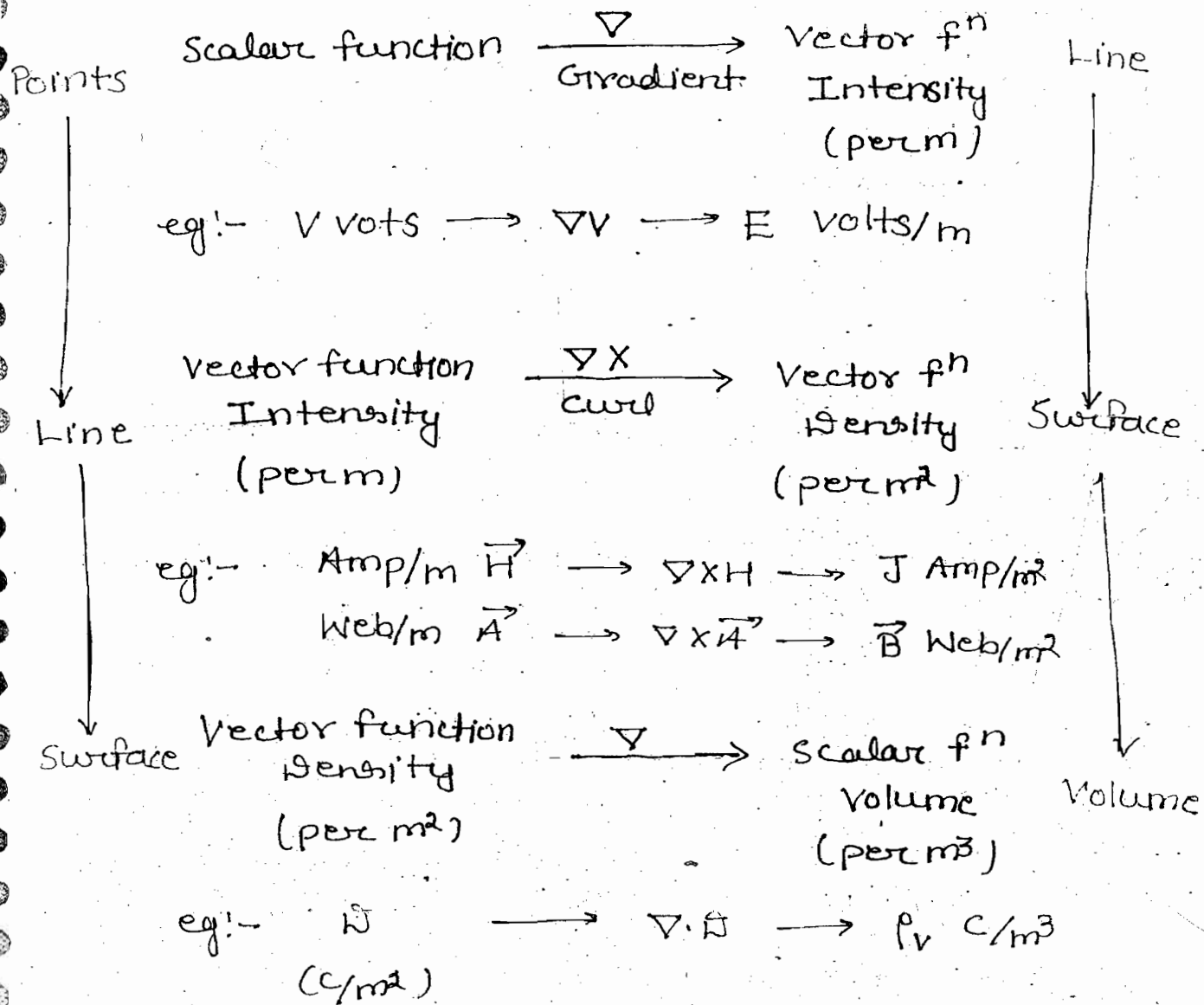
$$\vec{A} = \frac{W}{I d\vec{l}} = \frac{\mu I d\vec{l}}{4\pi r}$$

A's direction is along I flow direction & is a solenoidal vector.

Summary 1:-

- Q - Coulomb - point charge - Basic cause
(Scalar) - E field
- $I d\vec{l}$ - Amp-m - current element - H field
(Vector)
- $Q \rightarrow \rho_L d\vec{l} - \rho_S ds - \rho_V dv$
- $I d\vec{l} \rightarrow \vec{k} ds - \vec{J} dv$
- $\vec{E}$ - Intensity - Strength - Force - $\frac{\vec{F}}{Q} = \frac{1}{\epsilon}$
- $\vec{B}$ - Density - Strength - Force - $\frac{\vec{F}}{I d\vec{l}}$ dependent
- $\vec{E}$ - Density - Strength - $\frac{dQ}{ds} = C/m^2$ - μ dependent
- $\vec{H}$ - Intensity - Strength - $\frac{dI}{dl} = \text{Amp/m}$
- V - Potential - Strength - Work = $W/Q \rightarrow \text{Scalar}$
(Volts)
- $\vec{A}$ - Potential - Strength - Work = $\frac{W}{I d\vec{l}}$
(Weber/m) $\rightarrow \text{Vector}$

Summary :-



OR $\oint \vec{A} \cdot d\vec{l} = \frac{\text{Weber} \times \text{m}}{\text{m}} = \text{Webers}$ ↓ flux

$$\oint \vec{A} \cdot d\vec{l} = \int (\nabla \times \vec{A}) \cdot d\vec{s} = \int \vec{B} \cdot d\vec{s} = \psi_m$$

34. Stoke's theorem

$$\vec{V} = \nabla \times \vec{A}$$

C → closed line

S_C → surface

(B)

$$\oint \vec{A} \cdot d\vec{l} = \int (\nabla \times \vec{A}) \cdot d\vec{s} = \int \vec{V} \cdot d\vec{s}$$

Laplace / Poisson's Equations :-

→ They are second order differential equations relating volume charge and potential developed inside / outside the charge

$$\nabla \cdot \mathbf{D} = \rho_v$$

$$\nabla \cdot [\epsilon (-\nabla V)] = \rho_v$$

If $\epsilon = \text{constant}$ then

$$\boxed{\nabla^2 V = -\frac{\rho_v}{\epsilon}} \rightarrow \text{Poisson's Equation}$$

For $\rho_v = 0$, i.e. outside the charge, charge free region

$$\boxed{\nabla^2 V = 0} \rightarrow \text{Laplace Equation}$$

Note:-

In spite of being II order differential equation they always have a unique solution i.e. voltage is a single valued function of space

This is called as UNO Uniqueness theorem

Mathematically

$$\nabla^2 V = \nabla \cdot (\nabla V)$$

$$\nabla \cdot (\nabla V) = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u} \left(\frac{h_2 h_3}{h_1} \frac{\partial V}{\partial u} \right) + \frac{\partial}{\partial v} \left(\frac{h_3 h_1}{h_2} \frac{\partial V}{\partial v} \right) + \frac{\partial}{\partial w} \left(\frac{h_1 h_2}{h_3} \frac{\partial V}{\partial w} \right) \right]$$

Extension | —

$\vec{A}$ vs $\vec{J}$ in Magnetic Fields

$$\nabla \times \vec{H} = \vec{J}$$

$$\Rightarrow \nabla \times \left(\frac{\vec{B}}{\mu} \right) = \vec{J}$$

$$\Rightarrow \nabla \times (\nabla \times \vec{A}) = \mu \vec{J}$$

$$\Rightarrow \nabla (\nabla \cdot \vec{A}) - \nabla^2 \vec{A} = \mu \vec{J}$$

$$\Rightarrow \boxed{\nabla^2 \vec{A} = -\mu \vec{J}}$$

For current free region

$$\boxed{\nabla^2 \vec{A} = 0}$$

$$V = -\frac{6r^5}{\epsilon_0}$$

$$1. \nabla^2 V = -\rho_V / \epsilon$$

$$2. Q = \int \rho_V dV$$

$$\nabla^2 V = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \left(-\frac{6r^5}{\epsilon} \right) \right) = -\frac{\rho_V}{\epsilon}$$

$$\Rightarrow \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \cdot 6 \cdot 5 \cdot r^4) = \rho_V$$

$$\frac{30}{r^2} 6 \cdot r^5 = \rho_V$$

$$\Rightarrow \rho_V = 180 r^3$$

$$Q = \int_{r=0}^1 \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} 180 r^3 \cdot r^2 \sin \theta \, dr \, d\theta \, d\phi$$

$$= 180 \frac{\pi^6}{6} \Big|_0^1 (-\cos\theta)_0^\pi (\phi)_0^{2\pi} = 120\pi C, \text{ Ans.}$$

36.

$$\rho_v = -10^{-8} (1 + 10\rho)$$

$$V \text{ on } \varphi = 5\text{cm}$$

$$\text{Given } \rho = 2\text{cm} \quad V \text{ \& } E = 0$$

$$\nabla^2 V = -\frac{\rho_v}{\epsilon}$$

$$\nabla^2 V = \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \cdot \frac{\partial V}{\partial \rho} \right) = \frac{10^{-8} (1 + 10\rho)}{\frac{1}{36\pi \times 10^9}}$$

$$= 360\pi (1 + 10\rho)$$

$$\rho \frac{\partial V}{\partial \rho} = 360\pi \int (\rho + 10\rho^2) d\rho$$

$$\rho \frac{\partial V}{\partial \rho} = 360\pi \left(\frac{\rho^2}{2} + \frac{10\rho^3}{3} \right)$$

$$\Rightarrow V = 360\pi \int \left(\frac{\rho}{2} + \frac{10\rho^2}{3} \right) d\rho$$

$$= 360\pi \left(\frac{\rho^2}{4} + \frac{10\rho^3}{9} \right) \begin{matrix} 5 \times 10^{-2} \\ 2 \times 10^{-3} \end{matrix}$$

$$\Rightarrow \boxed{V = 0.5V}$$

37.

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} = -\frac{\rho_v}{\epsilon}$$

$$= 20 \cdot 3 \cdot 2 \cdot x + 10 \cdot 4 \cdot 3 \cdot y^2 \Big|_{(2,0)} = -\frac{\rho_v}{\epsilon}$$

$$\Rightarrow \boxed{\rho_v = -240\epsilon_0}$$

38. ϕ — 1 dimension function — satisfying — Laplace

$$\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} = 0$$

$$\frac{\partial \phi}{\partial x} = C_1 = \text{Rate of change is constant}$$

$$\frac{\phi_2 - \phi_1}{d} = \frac{\phi_3 - \phi_2}{2d}$$

$$\phi = C_1 x + C_2 \longrightarrow \text{Linear equation}$$

Ans — B

39.

$$V = \sinh x \cdot \cos ky \cdot e^{pz}$$

$$\nabla^2 V = 0$$

$$\frac{e^x + e^{-x}}{2} = \sinh x$$

$$\nabla^2 V = 1V - k^2 V + p^2 V = 0$$

$$\Rightarrow \boxed{k = \sqrt{1+p^2}}$$

Boundary Conditions for Electric fields

→ If a field is spread out into two different medium and known the field in one region. The field in the adjacent region can be calculated under two conditions

Case-(I) :-

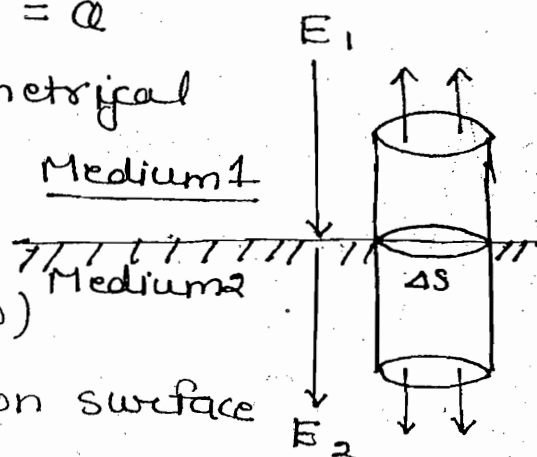
Electric field is normal to the boundary.

Using Gauss law $\oint \mathbf{E} \cdot d\mathbf{s} = Q$

consider a surface symmetrical in both the medium

$$\mathbf{E}_2 \cdot \Delta \mathbf{s} - \mathbf{E}_1 \cdot \Delta \mathbf{s} = 0$$

($\because Q=0$)



If charge is present on surface then

$$\mathbf{E}_2 \Delta \mathbf{s} - \mathbf{E}_1 \Delta \mathbf{s} = \rho_s \Delta \mathbf{s}$$

$$\mathbf{E}_{n1} = \mathbf{E}_{n2}$$

$$\mathbf{E}_{n2} - \mathbf{E}_{n1} = \rho_s$$

Statement:-

The normal components of electric flux density is same on either side but otherwise discontinuous when a surface charge density exists on the boundary

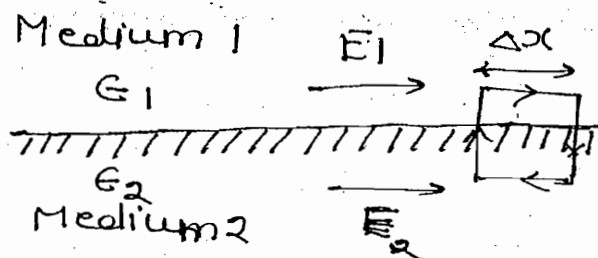
Case-(II) :-

Electric field is ^{tangential} parallel to boundary.

Using $\oint \mathbf{E} \cdot d\mathbf{l} = 0$

$$E_1 \Delta x - E_2 \Delta x = 0$$

$$\Rightarrow \boxed{E_1 = E_2}$$

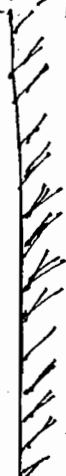


$$E_{t1} = E_{t2}$$

The tangential components of electric field intensity is always continuous

40

$x = 0$ Plane
 $YZ \rightarrow$ Plane

$x < 0$ $\epsilon_{R1} = 1.5$ $E_{n1} = 2a_x$ $E_{t1} = -3a_y + a_z$ $E_1 = 2a_x - 3a_y + a_z$		$x > 0$ $\epsilon_{R2} = 2.5$
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$$E_{t1} = E_{t2} = -3a_y + a_z$$

$$D_{n1} = 1.5 \times 2a_x \times \epsilon_0 \quad (D_{n1} = \epsilon \epsilon_r \times E_n)$$

$$= 3a_x = D_{n2}$$

$$= 3\epsilon_0 a_x \rightarrow D_{n2} = 3\epsilon_0 a_x$$

$$E_{n2} = \frac{3\epsilon_0 a_x}{2.5\epsilon_0} = 1.2a_x$$

$$D_1 = \epsilon_0 (3a_x - 4.5a_y + 1.5a_z)$$

$$D_2 = \epsilon_0 (3a_x - 7.5a_y + 2.5a_z)$$

$E_2 = 1.2a_x - 3a_y + a_z$

Ans

Note:-

→ The value of vector dec. then its projection dec. i.e. moves away from normal and moves towards the surface

→ The projection of normal component is dec. in E_2 when compare to E_1 . Hence the field is shifting away from the normal. It can be verified with inc. tangential component when Comparing D_1 & D_2

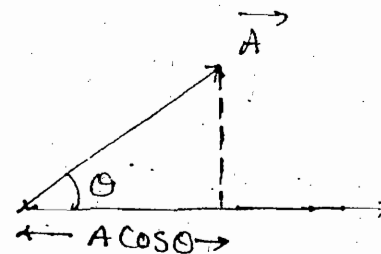
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A's scalar projection on B

$$= A \cos \theta$$

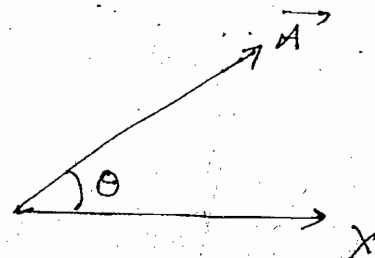
$$= \frac{A \cdot B}{|B|}$$

$$= \frac{A \cdot B}{|B|}$$



A's projection in x

$$= \frac{A \cdot a_x}{|a_x|}$$



A's vector projection on B

$$= (A \cos \theta) \cdot \hat{B}$$

$$= \frac{(A \cdot B)}{|B|} \hat{B}$$

$$= \frac{(A \cdot B)}{|B|^2} \vec{B}$$

41.

$$\rho_2 - \rho_1 = \rho_s$$

$$1 \cdot \epsilon_0 - 2 \cdot 2 \epsilon_0 = \rho_s$$

$$\rho_s = -3 \epsilon_0$$

$$E_2 = \epsilon_0$$

$$E_2 = q \alpha$$

$$E_1 = 2 \epsilon_0$$

$$E_1 = 2 q \alpha$$

42.

$$E_{t1} = E_{t2}$$

$$E_1 \cos 30^\circ = E_{t2}$$

$$\Rightarrow E_{t2} = 200 \sqrt{3}$$

$$\epsilon_0 E_1 \sin 30^\circ = 20 \epsilon_0 E_{n2}$$

$$\Rightarrow E_{n2} = 10$$

