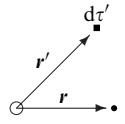


## 7.3 Electromagnetic fields (general)

### Field relationships

|   |   |        |   |
|---|---|--------|---|
| Conservation of charge  | $\nabla \cdot \mathbf{J} = -\frac{\partial \rho}{\partial t}$   | (7.39) | $\mathbf{J}$ current density<br>$\rho$ charge density<br>$t$ time                   |
| Magnetic vector potential   | $\mathbf{B} = \nabla \times \mathbf{A}$   | (7.40) | $\mathbf{A}$ vector potential   |
| Electric field from potentials                                    | $\mathbf{E} = -\frac{\partial \mathbf{A}}{\partial t} - \nabla \phi$  | (7.41) | $\phi$ electrical potential   |
| Coulomb gauge condition   | $\nabla \cdot \mathbf{A} = 0$   | (7.42) |   |
| Lorenz gauge condition  | $\nabla \cdot \mathbf{A} + \frac{1}{c^2} \frac{\partial \phi}{\partial t} = 0$  | (7.43) | $c$ speed of light  |
| Potential field equations <sup>a</sup>                            | $\frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} - \nabla^2 \phi = \frac{\rho}{\epsilon_0}$  | (7.44) |  |
|   | $\frac{1}{c^2} \frac{\partial^2 \mathbf{A}}{\partial t^2} - \nabla^2 \mathbf{A} = \mu_0 \mathbf{J}$   | (7.45) |   |
| Expression for $\phi$ in terms of $\rho$ <sup>a</sup>             | $\phi(\mathbf{r}, t) = \frac{1}{4\pi\epsilon_0} \int_{\text{volume}} \frac{\rho(\mathbf{r}', t -  \mathbf{r} - \mathbf{r}' /c)}{ \mathbf{r} - \mathbf{r}' } d\tau'$       | (7.46) | $d\tau'$ volume element<br>$\mathbf{r}'$ position vector of $d\tau'$                |
| Expression for $\mathbf{A}$ in terms of $\mathbf{J}$ <sup>a</sup> | $\mathbf{A}(\mathbf{r}, t) = \frac{\mu_0}{4\pi} \int_{\text{volume}} \frac{\mathbf{J}(\mathbf{r}', t -  \mathbf{r} - \mathbf{r}' /c)}{ \mathbf{r} - \mathbf{r}' } d\tau'$ | (7.47) | $\mu_0$ permeability of free space  |

<sup>a</sup> Assumes the Lorenz gauge.

### Liénard–Wiechert potentials<sup>a</sup>

|  |  |        |   |
|--|--|--------|---|
| Electrical potential of a moving point charge      | $\phi = \frac{q}{4\pi\epsilon_0( \mathbf{r}  - \mathbf{v} \cdot \mathbf{r}/c)}$              | (7.48) | $q$ charge<br>$\mathbf{r}$ vector from charge to point of observation<br>$\mathbf{v}$ particle velocity |
| Magnetic vector potential of a moving point charge | $\mathbf{A} = \frac{\mu_0 q \mathbf{v}}{4\pi( \mathbf{r}  - \mathbf{v} \cdot \mathbf{r}/c)}$ | (7.49) |                     |

<sup>a</sup>In free space. The right-hand sides of these equations are evaluated at retarded times, i.e., at  $t' = t - |\mathbf{r}'|/c$ , where  $\mathbf{r}'$  is the vector from the charge to the observation point at time  $t'$ .

## Maxwell's equations

|  |   |
|--|---|
| Differential form:   | Integral form:  |
| $\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$ (7.50)   | $\oint_{\text{closed surface}} \mathbf{E} \cdot d\mathbf{s} = \frac{1}{\epsilon_0} \int_{\text{volume}} \rho d\tau$ (7.51)  |
| $\nabla \cdot \mathbf{B} = 0$ (7.52)   | $\oint_{\text{closed surface}} \mathbf{B} \cdot d\mathbf{s} = 0$ (7.53)   |
| $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$ (7.54)                                    | $\oint_{\text{loop}} \mathbf{E} \cdot d\mathbf{l} = -\frac{d\Phi}{dt}$ (7.55)   |
| $\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$ (7.56) | $\oint_{\text{loop}} \mathbf{B} \cdot d\mathbf{l} = \mu_0 I + \mu_0 \epsilon_0 \int_{\text{surface}} \frac{\partial \mathbf{E}}{\partial t} \cdot d\mathbf{s}$ (7.57) |
| Equation (7.51) is “Gauss’s law”   | $d\mathbf{s}$ surface element   |
| Equation (7.55) is “Faraday’s law”   | $d\tau$ volume element  |
| $\mathbf{E}$ electric field  | $d\mathbf{l}$ line element  |
| $\mathbf{B}$ magnetic flux density   | $\Phi$ linked magnetic flux ( $= \int \mathbf{B} \cdot d\mathbf{s}$ )   |
| $\mathbf{J}$ current density   | $I$ linked current ( $= \int \mathbf{J} \cdot d\mathbf{s}$ )  |
| $\rho$ charge density  | $t$ time  |

## Maxwell's equations (using $\mathbf{D}$ and $\mathbf{H}$ )

|   |  |
|---|--|
| Differential form:  | Integral form:   |
| $\nabla \cdot \mathbf{D} = \rho_{\text{free}}$ (7.58)   | $\oint_{\text{closed surface}} \mathbf{D} \cdot d\mathbf{s} = \int_{\text{volume}} \rho_{\text{free}} d\tau$ (7.59)  |
| $\nabla \cdot \mathbf{B} = 0$ (7.60)  | $\oint_{\text{closed surface}} \mathbf{B} \cdot d\mathbf{s} = 0$ (7.61)  |
| $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$ (7.62)   | $\oint_{\text{loop}} \mathbf{E} \cdot d\mathbf{l} = -\frac{d\Phi}{dt}$ (7.63)  |
| $\nabla \times \mathbf{H} = \mathbf{J}_{\text{free}} + \frac{\partial \mathbf{D}}{\partial t}$ (7.64)                                       | $\oint_{\text{loop}} \mathbf{H} \cdot d\mathbf{l} = I_{\text{free}} + \int_{\text{surface}} \frac{\partial \mathbf{D}}{\partial t} \cdot d\mathbf{s}$ (7.65) |
| $\mathbf{D}$ displacement field   | $\mathbf{E}$ electric field  |
| $\rho_{\text{free}}$ free charge density (in the sense of<br>$\rho = \rho_{\text{induced}} + \rho_{\text{free}}$ )                          | $d\mathbf{s}$ surface element  |
| $\mathbf{B}$ magnetic flux density  | $d\tau$ volume element   |
| $\mathbf{H}$ magnetic field strength  | $d\mathbf{l}$ line element   |
| $\mathbf{J}_{\text{free}}$ free current density (in the sense of<br>$\mathbf{J} = \mathbf{J}_{\text{induced}} + \mathbf{J}_{\text{free}}$ ) | $\Phi$ linked magnetic flux ( $= \int \mathbf{B} \cdot d\mathbf{s}$ )  |
|   | $I_{\text{free}}$ linked free current ( $= \int \mathbf{J}_{\text{free}} \cdot d\mathbf{s}$ )  |
|   | $t$ time   |

## Relativistic electrodynamics

|  |   |        |   |
|--|---|--------|---|
| Lorentz transformation of electric and magnetic fields | $E'_\parallel = E_\parallel$  | (7.66) | $\mathbf{E}$ electric field                         |
|  | $E'_\perp = \gamma(\mathbf{E} + \mathbf{v} \times \mathbf{B})_\perp$                  | (7.67) | $\mathbf{B}$ magnetic flux density                  |
|  | $\mathbf{B}'_\parallel = \mathbf{B}_\parallel$  | (7.68) | ' measured in frame moving at relative velocity $v$ |
|  | $\mathbf{B}'_\perp = \gamma(\mathbf{B} - \mathbf{v} \times \mathbf{E}/c^2)_\perp$     | (7.69) | $\gamma$ Lorentz factor<br>= $[1 - (v/c)^2]^{-1/2}$ |
| Lorentz transformation of current and charge densities | $\rho' = \gamma(\rho - v J_\parallel/c^2)$  | (7.70) | $\parallel$ parallel to $v$                         |
|  | $J'_\perp = J_\perp$  | (7.71) | $\perp$ perpendicular to $v$                        |
|  | $J'_\parallel = \gamma(J_\parallel - v\rho)$  | (7.72) |   |
| Lorentz transformation of potential fields             | $\phi' = \gamma(\phi - v A_\parallel)$  | (7.73) | $\mathbf{J}$ current density                        |
|  | $A'_\perp = A_\perp$  | (7.74) | $\rho$ charge density                               |
|  | $A'_\parallel = \gamma(A_\parallel - v\phi/c^2)$                                      | (7.75) | $\phi$ electric potential                           |
|  | $\tilde{\mathbf{J}} = (\rho c, \mathbf{J})$   | (7.76) | $\mathbf{A}$ magnetic vector potential              |
| Four-vector fields <sup>a</sup>                        | $\tilde{\mathbf{A}} = \left( \frac{\phi}{c}, \mathbf{A} \right)$                      | (7.77) |   |
|  | $\square^2 = \left( \frac{1}{c^2} \frac{\partial^2}{\partial t^2}, -\nabla^2 \right)$ | (7.78) | $\tilde{\mathbf{J}}$ current density four-vector    |
|  | $\square^2 \tilde{\mathbf{A}} = \mu_0 \tilde{\mathbf{J}}$                             | (7.79) | $\tilde{\mathbf{A}}$ potential four-vector          |
|  |   |        | $\square^2$ D'Alembertian operator                  |

<sup>a</sup>Other sign conventions are common here. See page 65 for a general definition of four-vectors.